

## Green's functions and propagators

Useful relations :

$$\int_{-\infty}^{+\infty} dy e^{iyz} = 2\pi\delta(z); \quad \frac{d}{dz}\theta(z) = \delta(z)$$

Cauchy's theorem :

$$\oint dz \frac{f(z)}{z-w} = 2n\pi i f(w)$$

where  $n$  is the number of times the integration contour goes around  $w$  anti-clockwise.

### 1. KG equation coupled to an external source

Consider the equation for a real scalar field with an external source :

$$(\square + m^2)\phi(\vec{x}, t) = J(\vec{x}, t). \tag{1}$$

1. Assuming there is an instant  $t_0$  such that  $J(\vec{x}, t < t_0) = 0$ , write, in terms of an appropriate Green's function, the solution  $\phi(\vec{x}, t)$  which satisfies  $\phi(\vec{x}, t < t_0) = 0$ . Which is the correctn Green's function in this case?
2. Find the explicit solution  $\phi(\vec{x}, t)$  in the particular case :

$$J(\vec{x}, t) = j_0 \theta(t) e^{-\mu t}, \quad j_0, \mu > 0, \quad \theta(t) = \begin{cases} 1 & t \geq 0 \\ 0 & t < 0 \end{cases} \tag{2}$$

3. Recall the solution for the classical Klein-Gordon (1) in the absence of sources :

$$\phi(\vec{x}, t) = \int \frac{d^3k}{(2\pi)^3 \sqrt{2\omega_k}} \left( a(\vec{k}) e^{i\vec{k}\cdot\vec{x} - i\omega_k t} + a^*(\vec{k}) e^{-i\vec{k}\cdot\vec{x} + i\omega_k t} \right) \tag{3}$$

where  $\omega_k \equiv \sqrt{\vec{k}^2 + m^2}$ . Show that, for  $t \rightarrow \infty$ , the solution found in the previous point reduces to a solution of the free equation with coefficients :

$$a(\vec{k}) = -\frac{(2\pi)^{3/2} j_0}{(m + i\mu)\sqrt{2m}} \delta^3(\vec{k}). \tag{4}$$

### 3. Green's function for Helmholtz equation

Consider the following static equation in  $d=3$  space dimensions (no time) :

$$(-\nabla^2 + m^2)\phi(\vec{x}) = 0 \tag{5}$$

where  $\phi(\vec{x})$  is a scalar field (with respect rotations in 3d). We want to construct the associated Green's function,  $G(\vec{x})$ , i.e. the solution of

$$(-\nabla^2 + m^2)G(\vec{x}) = \delta^{(3)}(\vec{x}). \tag{6}$$

1. Writing equation (6) in Fourier space, show that :

$$G(\vec{x}) = \int \frac{d^3k}{(2\pi)^3} \frac{e^{i\vec{k}\cdot\vec{x}}}{k^2 + m^2}. \tag{7}$$

Is the Green's function unique? Do we need to go through the contour-deforming procedure in this case? Why is that the case?

2. Perform the 3d momentum integral explicitly and show that :

$$G(\vec{x}) = \frac{e^{-m|\vec{x}|}}{4\pi|\vec{x}|} \quad (8)$$

[Hint : choose spherical coordinate \*wisely\*, then perform the integral over angles, then use a contour integral for the final integration over  $|\vec{k}|$ . ]

#### 4. The propagator as a Green function

1. Using the definition of the Feynman, retarded and advanced contours, show that :

$$G_F(x) = G_A(x) + iD_+(x), \quad \text{and} \quad G_F(x) = G_R(x) + iD_-(x) \quad (9)$$

where  $G_A$  and  $G_R$  are the advanced and retarded Green's functions, and  $D_{\pm}$  are the positive and negative frequency *Wightman functions*, defined by :

$$D_{\pm}(x) = \int \frac{d^3p}{(2\pi)^3} \frac{e^{\mp i(\omega_p t - \vec{p} \cdot \vec{x})}}{2\omega_p}$$

where  $\omega_p = \sqrt{|\vec{p}|^2 + m^2}$ . Verify that  $D_+(x) = D_-(-x)$ .

2. Using the results above, show that

$$G_F(x - x') = i \left( \theta(t - t') D_+(x - x') + \theta(t' - t) D_-(x - x') \right). \quad (10)$$

3. By explicit calculation applying the differential KG operator in position space show that the right hand side of equation (10) is a Green's function for the Klein-Gordon equation.
4. In free Klein-Gordon theory, compute the two-point function :

$$\langle 0 | \hat{\phi}(x) \hat{\phi}(x') | 0 \rangle,$$

and give its relation to  $D_+(x - x')$  and  $D_-(x - x')$ . Use this to obtain the result :

$$G_F(x - x') = i \langle 0 | T \left( \hat{\phi}(x) \hat{\phi}(x') \right) | 0 \rangle.$$

where  $T$  denotes time-ordering.

This gives the fundamental result that the Feynman Green's function coincides with the time-ordered two-point correlator.

5. Similarly, show that

$$G_R(x, x') = i \theta(t - t') \langle 0 | \left[ \hat{\phi}(x), \hat{\phi}(x') \right] | 0 \rangle. \quad (11)$$

#### 5. Lorentz-invariant measure

Show that the integration measure

$$\frac{d^3p}{2\omega_p} \quad (12)$$

is Lorentz-invariant (where  $\omega_p \equiv \sqrt{|\vec{p}|^2 + m^2}$ ).

[Hint : rewrite (12) as an integration measure over 4-momentum with suitable Dirac delta-functions]