

From Nuclei to stars: The nucleosynthesis processes at work in Universe: from Big-Bang to stars

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Plan of lecture II

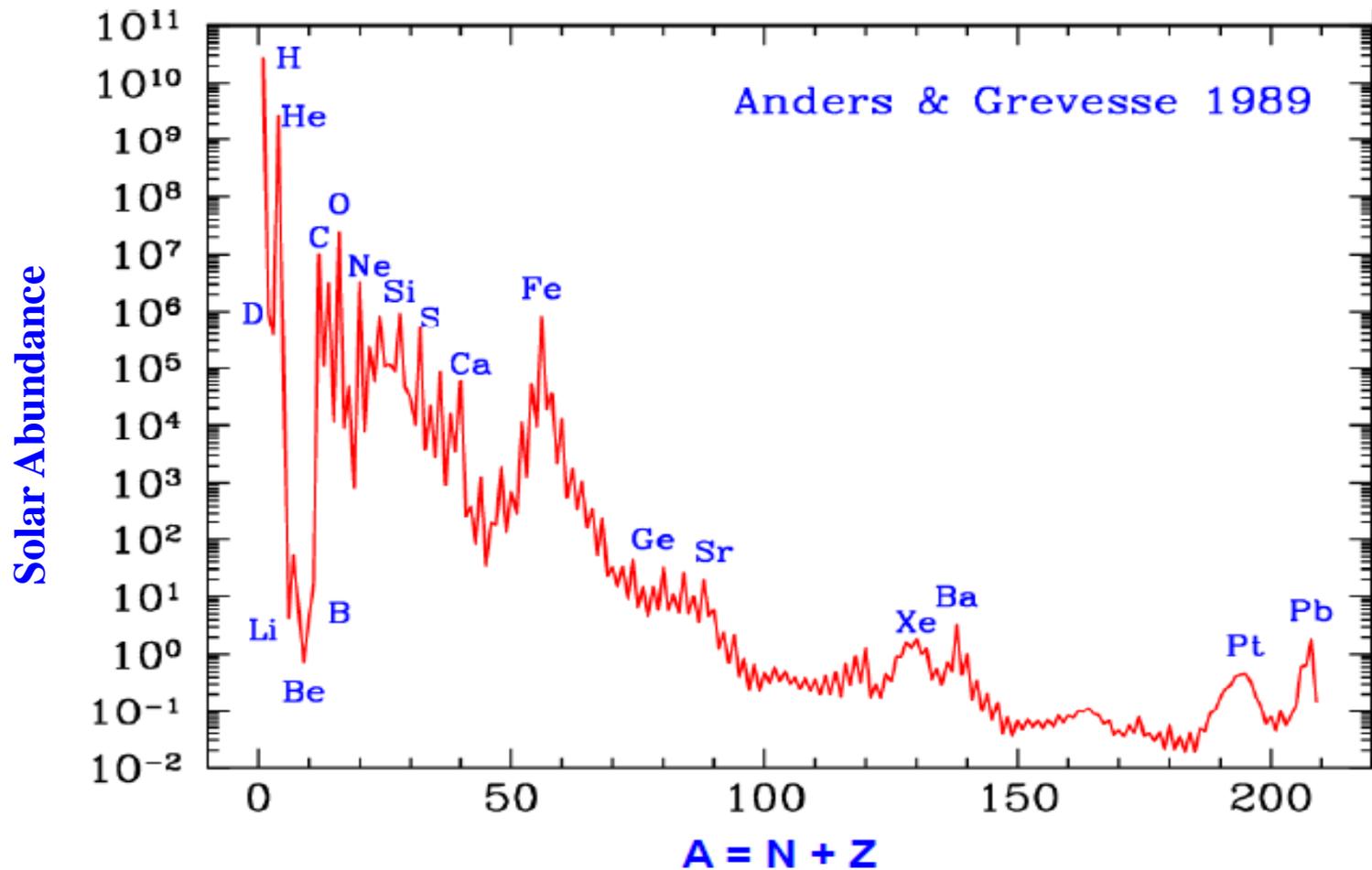
- I. Little bit of history
- II. Big-Bang nucleosynthesis
- III. Cosmic ray nucleosynthesis
- IV. Stellar nucleosynthesis
 - Hydrogen burning: p-p chains and CNO cycles
 - Helium burning
 - Advanced burning stages: C-, Ne, O and Si burning
 - Explosive burning in Core Collapse Supernovae
 - Explosive nucleosynthesis
 - Nucleosynthesis beyond iron: s- and r-process

Text books

- **Cauldrons in the Cosmos**, Nuclear Astrophysics ,
Claus E. Rolfs and William S.Rodney
The University of Chicago Press, 1988
ISBN 0-226-72456-5
- **Principles of Stellar Evolution and Nucleosynthesis**,
Donald D. Clayton, The University of Chicago Press ,1968
ISBN 0-226-10953-4
- **Supernovae and Nucleosynthesis** – An investigation of the History of
Matter, from the Big-Bang to the Present.
David Arnett

Nucleosynthesis : When and where?

Abundance curve of the elements:



- 1920 Eddington suggests that the energy of the stars result from nuclear fusion
- 1928 Gamow : discovery of quantum tunneling, which demonstrates the non-zero probability of nuclear reactions in stars
- 1938 Bethe, Critchfield & von Weizsaecker : discovery of pp chain and CNO cycle
- 1951–1953 Öpik, Salpeter and Hoyle : discovery of He burning mechanism
- 1957 Burbidge, Burbidge, Fowler & Hoyle (B²FH) give an overview of nucleosynthesis processes in the Stars
- 1968 Davis et al. : 1st detection of neutrinos emitted in Sun core
- 1987 γ radioactivity detection of ⁵⁶Co and ⁵⁷Co in supernova SN 1987A
- 2013 Observational evidence of heavy nuclei nucleosynthesis in the coalescence of a binary system of two neutron stars (GRB 130603B)
- **2017 (August)** observational confirmation of heavy nuclei nucleosynthesis in a binary neutron star mergers (GW170817)

Primordial Nucleosynthesis



Alpher



Bethe
("α β γ")



Gamow

“All the elements were formed
just after Big-Bang”

Phys. Rev. 73. (1948) 803

→ Almost right for D, He, a part of ${}^7\text{Li}$
But: Mass gap $A=5$ & $A=8$ (unbound nuclei)

Stellar Nucleosynthesis



Burbidge



Burbidge
(B2FH)



Fowler



Hoyle



REVIEWS OF

MODERN PHYSICS

VOLUME 29, NUMBER 4

OCTOBER, 1957

Synthesis of the Elements in Stars*

E. MARGARET BURBIDGE, G. R. BURBIDGE, WILLIAM A. FOWLER, AND F. HOYLE

*Kellogg Radiation Laboratory, California Institute of Technology, and
Mount Wilson and Palomar Observatories, Carnegie Institution of Washington,
California Institute of Technology, Pasadena, California*

“All elements are synthesized in stars
through various processes”

Rev. Mod. Phys. 29 (1957) 547

B2FH epoch:

- Primordial (Big-Bang) nucleosynthesis
- H burning & He burning
- “e” process (iron peak)
- “x” process (LiBeB, “x” for unknown)
- “r” process (rapid neutron capture)
- “s” process (slow neutron capture)
- “p” process (proton rich)

Today : no fundamental changes, “x” → identified as non-thermal nucleosynthesis and C,Ne, O, Si burning phases added to the list above.

Big-Bang Nucleosynthesis

Three observational pillars of Big-Bang Model:

➤ The expansion of the Universe

→ Galaxies move away from each other & from us according to Hubble's law: $V = H_0 \times D$ with $H_0 \approx 70 \text{ km/s/Mpc}$, the Hubble constant

➤ Cosmic Microwave Background radiation (CMB)

→ A black body radiation at **2.7 K** corresponding to the redshifted spectrum emitted when the universe became transparent (Penzias & Wilson in 1965)

➤ Primordial nucleosynthesis (BBN) of light elements

→ BBN reproduces the observed primordial abundances over a range of **nine order of magnitudes**

➤ At $T > 10^{10}$ K, the particles present: photons, e^- , e^+ , ν , anti- ν , p , n

Equilibrium $p \leftrightarrow n$: $N_n/N_p = \exp(-Q_{np}/kT)$; $Q_{np} = 1.29 \text{ MeV}$



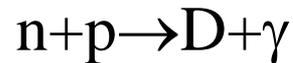
Equilibrium as long as the weak reaction rate is faster than the expansion rate: $\Gamma_{n \leftrightarrow p} \gg H(t)$

Equilibrium breaks out when :

$$\Gamma_{n \leftrightarrow p} < H(t) \quad \text{at } T \approx 10^{10} \text{ K}$$

Then $N_n/N_p \simeq 0.13$ due to free neutron β decay

Neutrons decay until T is low enough for :



becomes faster than deuterium photodisintegration



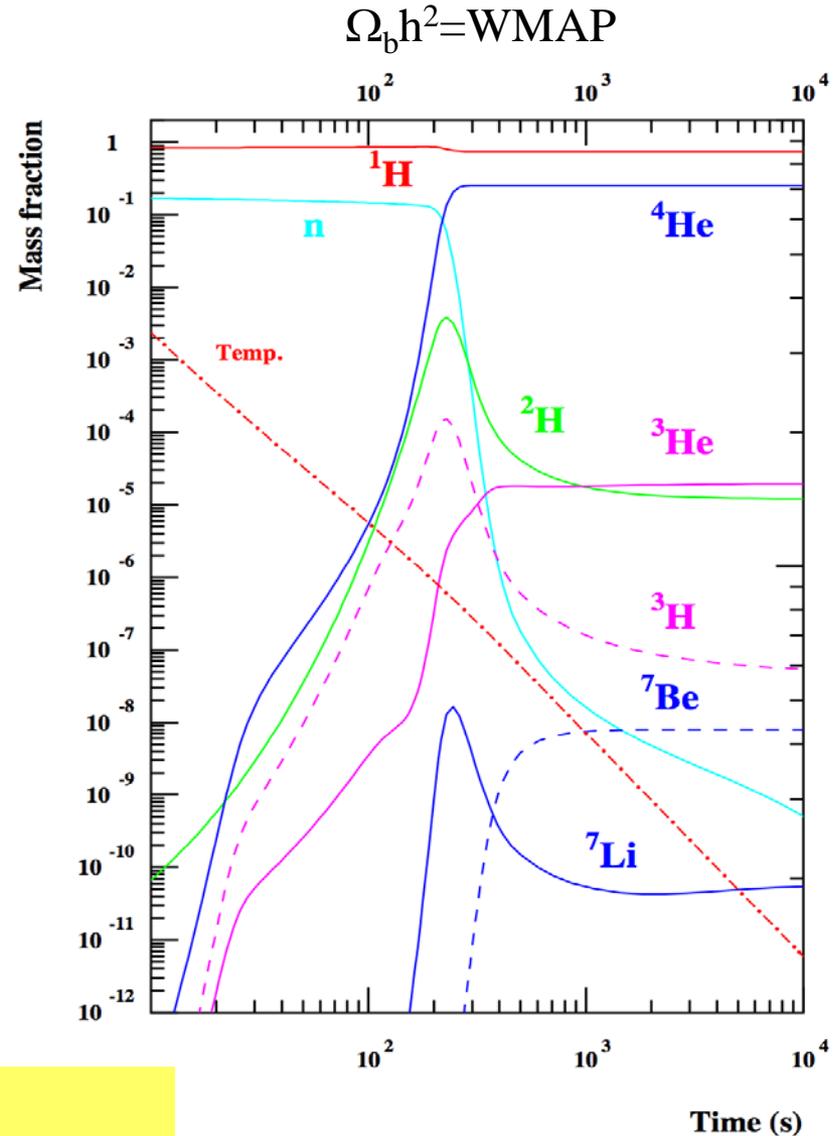
Then, $t = 3 \text{ mn}$, $T \approx 10^9 \text{ K}$ and N_n has decreased to $N_n/N_p \approx 0.1$

Nucleosynthesis starts to produce essentially ${}^4\text{He}$ together with traces of D , ${}^3\text{He}$, ${}^7\text{Li}$,

Big-Bang nucleosynthesis: The canonical BBN reaction network

The 12 reactions of standard BBN:

- $n \leftrightarrow p$ with $\tau_n = 880 \pm 4$ s
- $p + n \rightarrow D + \gamma$
- $D + p \rightarrow {}^3\text{He} + \gamma$
- $D + D \rightarrow {}^3\text{He} + n$
- $D + D \rightarrow T + p$
- $T + D \rightarrow {}^4\text{He} + n$
- $T + {}^4\text{He} \rightarrow {}^7\text{Li} + \gamma$
- ${}^3\text{He} + n \rightarrow p + T$
- ${}^3\text{He} + D \rightarrow p + {}^4\text{He}$
- ${}^3\text{He} + {}^4\text{He} \rightarrow {}^7\text{Be} + \gamma$
- ${}^7\text{Li} + p \rightarrow {}^4\text{He} + {}^4\text{He}$
- ${}^7\text{Be} + n \rightarrow {}^7\text{Li} + p$



Number of baryons per photon : $\eta \equiv n_b/n_\gamma$

Baryonic density of Universe $\Omega_b h^2 = 3.65 \times 10^7 \eta$

Big-Bang nucleosynthesis:

Observations versus predictions

Observations: From a set of the most primitive astrophysical objects with the lowest metallicity

- **D observations** in clouds (cold objects) at high redshift on the line of sight of distant quasars

$$D/H = (2.53 \pm 0.40) \times 10^{-5}$$

- **^4He observations** in HII (ionized H) regions of blue compact galaxies

$$Y_p = 0.2561 \pm 0.0108$$

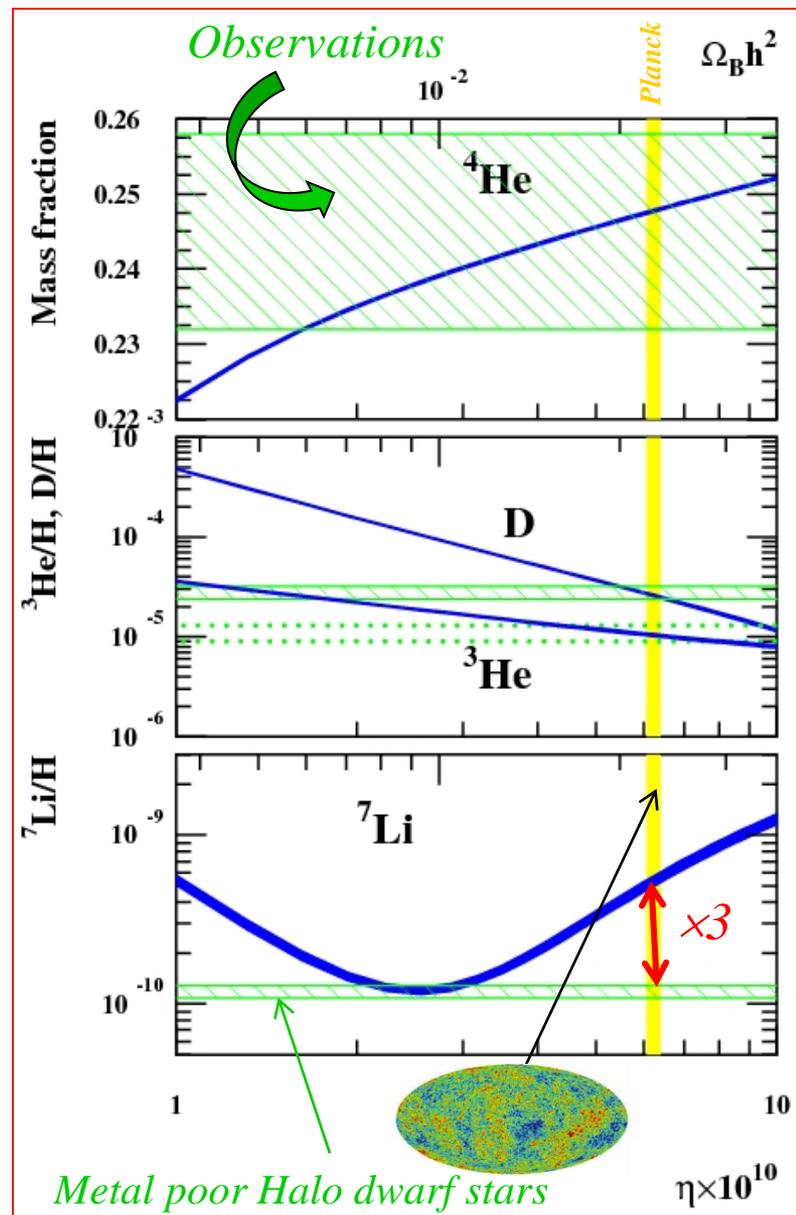
- **^3He observations** in HII regions of our Galaxy

$$^3\text{He}/H = (1.1 \pm 0.2) \times 10^{-5}$$

- **^7Li observations** at the surface of metal poor Halo dwarf stars

$$^7\text{Li}/H = (1.58 \pm 0.31) \times 10^{-10} = 1/3 \text{ predicted } ^7\text{Li}$$

→ The ^7Li cosmological problem



What are the possible solutions to ${}^7\text{Li}$ problem?

- **Astrophysical** solution?
- **Nuclear physics** solution?
- **Physics beyond** standard model?

Big-Bang nucleosynthesis: The astrophysical solution (I) to ${}^7\text{Li}$?

→ Primordial ${}^7\text{Li}$ abundance measured in old metal poor halo dwarf stars

Spite plateau:

Li/H versus [Fe/H] (time)

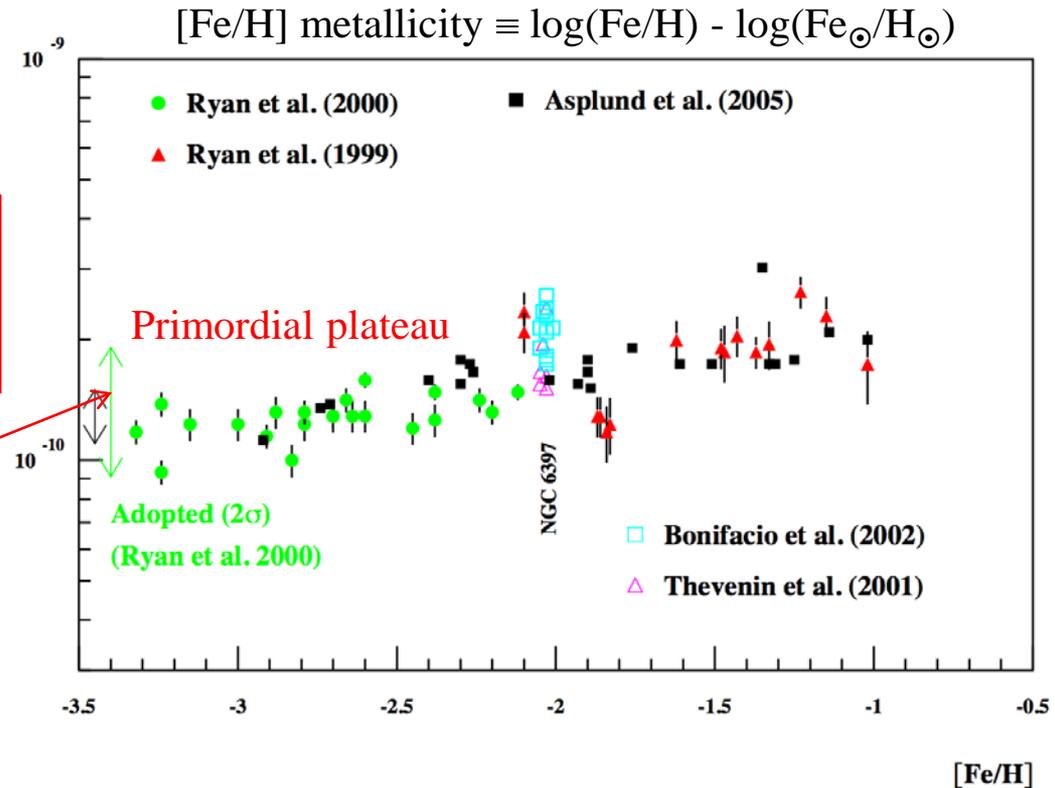
Li/H $\approx 1.12 \times 10^{-10}$ [Spite & Spite, 1982]

$0.9 \leftrightarrow 1.9 \times 10^{-10}$ [Ryan et al. (2000)]

$\approx 1.3 \leftrightarrow 2.3 \times 10^{-10}$ [Charbonnel et al (2005)]

$\approx 1.1 \leftrightarrow 1.5 \times 10^{-10}$ [Asplund et al. (2006)]

- Very low dispersion



➤ Reliable primordial abundance?

→ Systematic errors in the extraction of Li abundance due to the atmosphere models ?

Big-Bang nucleosynthesis: The astrophysical solution (II) to ${}^7\text{Li}$?

→ **Depletion** of the **atmospheric ${}^7\text{Li}$** due to rotationally induced **mixing and/diffusion?**

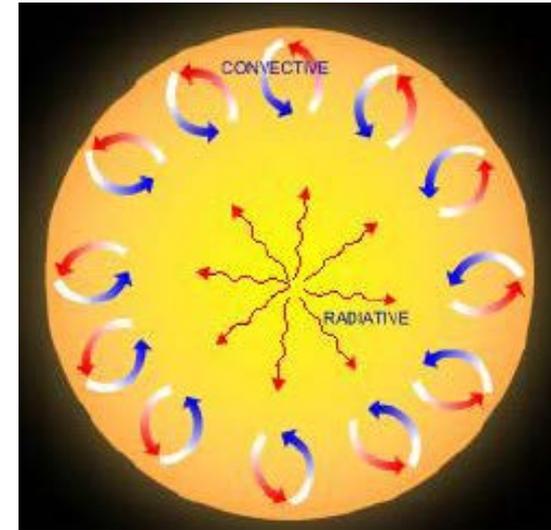
➤ Lithium easily burned in stars (**low binding energy**):
 $T \geq 2.5 \times 10^6 \text{ K} \Rightarrow {}^7\text{Li} + \text{p} \rightarrow 2 {}^4\text{He}$

➤ Destroyed by convection

- Deeper mixing of surface material → ${}^7\text{Li}$ burning

➤ **But not enough and not uniformity!**

- **Metal poor** halo stars have **shallower convective zones** than in solar metallicity stars
- Different stars have different convective zones
- Should see **larger scatter** around “**plateau**”



Big-Bang nucleosynthesis:

Nuclear solution to ${}^7\text{Li}$ problem?

➤ ${}^7\text{Li}$ production via **EC** decay of ${}^7\text{Be}$ at high η :

➤ Main ${}^7\text{Be}$ production mechanism: ${}^3\text{He}({}^4\text{He},\gamma){}^7\text{Be}$

→ Various measurements of the cross-section:
10% uncertainty

➤ Main ${}^7\text{Be}$ destruction mechanism: ${}^7\text{Be}(n,p){}^7\text{Li}(p,\alpha)\alpha$

→ ${}^7\text{Be}(n,p){}^7\text{Li}$ cross-section well known :

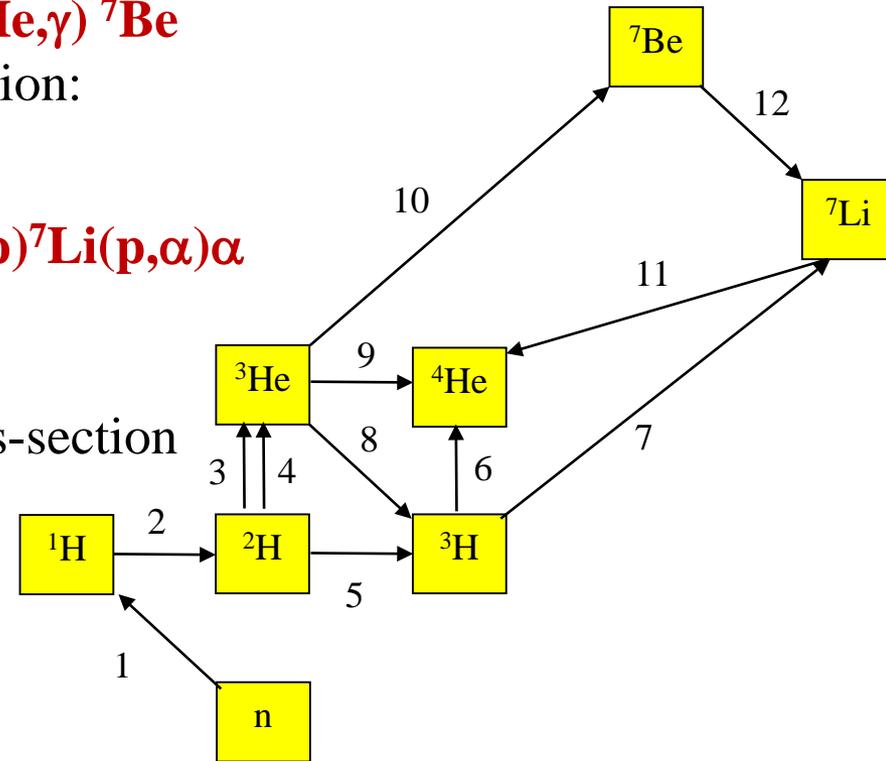
1% uncertainty

→ ${}^7\text{Li}(p,\alpha)\alpha$: 6% uncertainty on the cross-section

➤ Secondary destruction reactions:

${}^7\text{Be}+d$, ${}^7\text{Be}+{}^3\text{He}$, ${}^7\text{Be}+{}^4\text{He}$, ...

→ All experimentally studied these last years & **none is responsible** of the ${}^7\text{Li}$ deficit



The solution to the ${}^7\text{Li}$ problem has very likely to be found outside nuclear physics

→ Decays of heavier meta-stable (100-1000s) particles that **inject** additional **neutrons** which could increase the rates of ${}^7\text{Be}+n$

[Jedamzik (2004,2006), Kawasaki et al. (2005), Ellis et al. (2005)]

→ Existence of **mirror universe** in which neutrons can oscillate to our world \Rightarrow effective late time **neutron injection** which induces an increase of the ${}^7\text{Be}$ **destruction** by ${}^7\text{Be}+n$ [Coc et al. 2013]

BBN produced **H, D, He** & part of **Li** only. The nucleosynthesis of heavier element is prevented by:

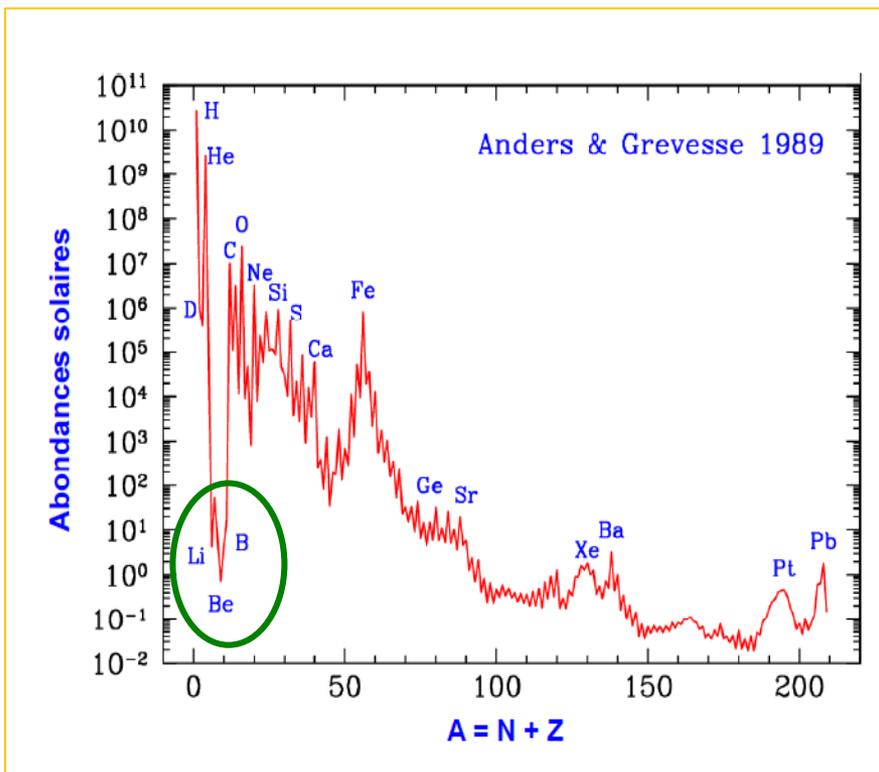
- coulomb barriers (see next course) of elements with higher atomic numbers
- the lack of stable isotopes of mass numbers **5** and **8**
- the decreasing density of matter as the universe expands

BBN details => see cosmology course

Cosmic-ray Nucleosynthesis

Cosmic ray nucleosynthesis:

The LiBeB story

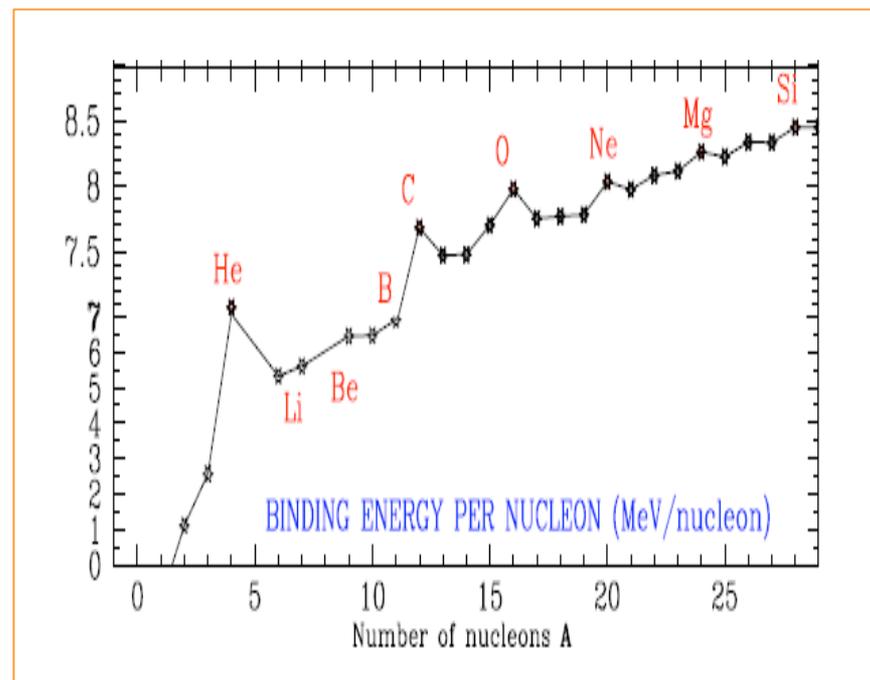


➤ BBN model: 10 order magnitudes differences between ${}^7\text{Li}$ and ${}^{10}\text{B}$

⇒ Light elements **not from BBN**

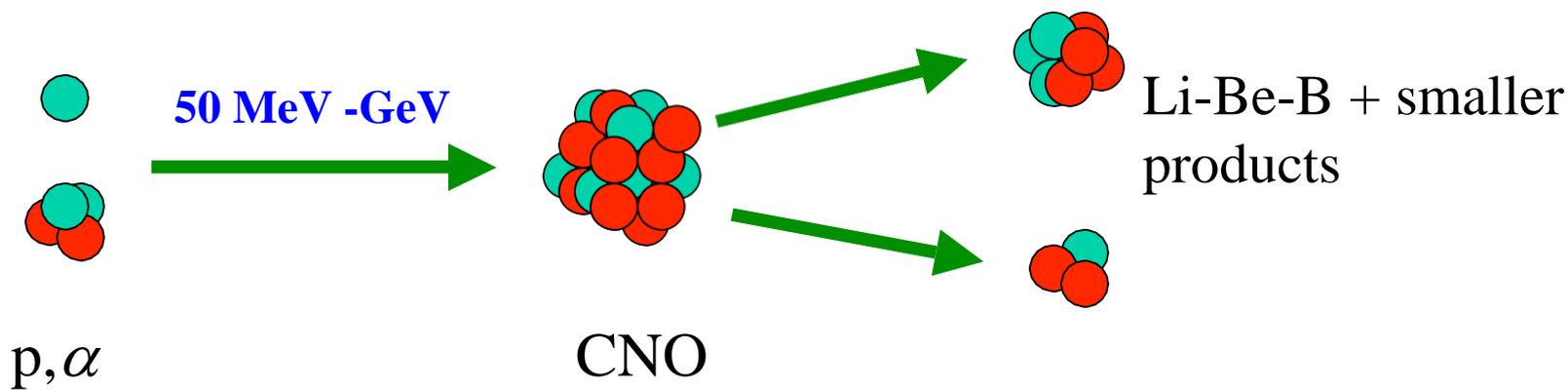
➤ ${}^6\text{Li}$, ${}^7\text{Li}$, ${}^9\text{Be}$, ${}^{10}\text{B}$, ${}^{11}\text{B}$ are fragile enough to be destroyed in stellar interiors

⇒ **Not in stars**



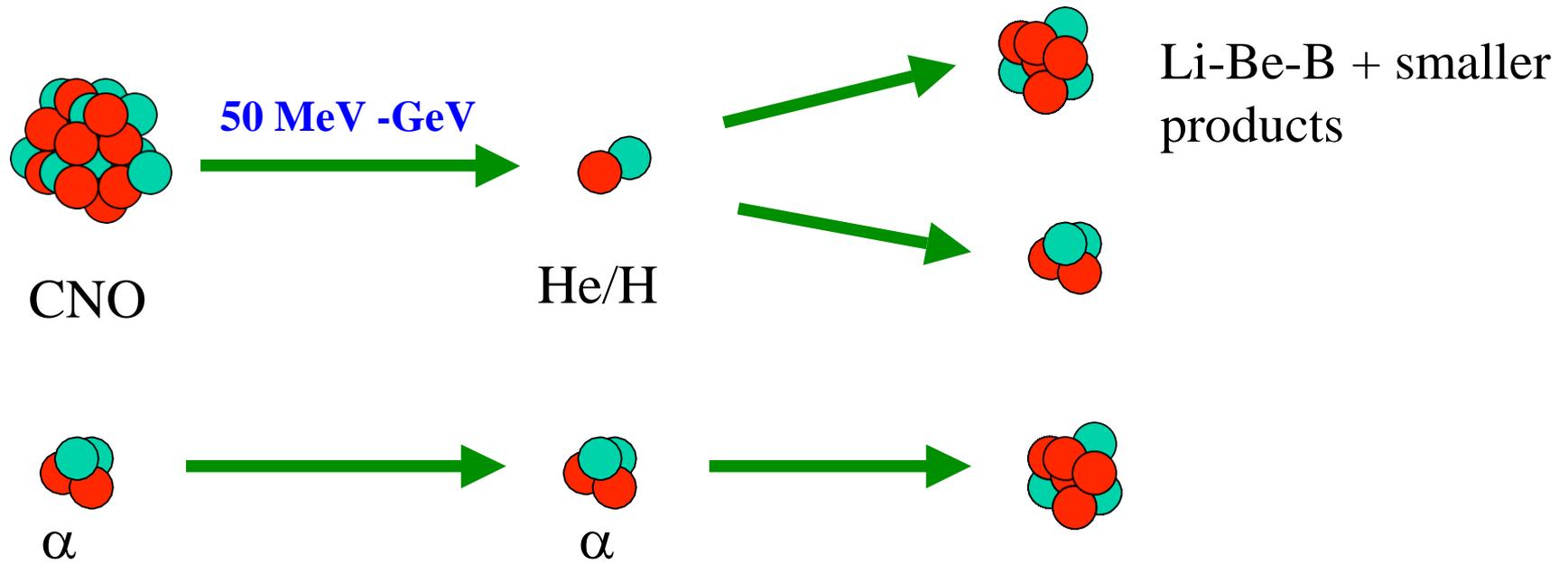
➤ **Must come from galactic processes: Cosmic ray nucleosynthesis**

→ **Spallation**



- Cosmic ray particles with high energy (50 MeV – GeV) interact with CNO in ISM to produce light elements
- This is a *metallicity dependent* “secondary” process – it depends on the amount of CNO in the ISM

→ Inverse Spallation



- C, N, O, H accelerated by supernovae reacts with interstellar H & He, breaks apart

**Stellar Nucleosynthesis:
Hydrostatic (quiescent)
stellar burning**

Little Reminder

To achieve life as a star for a protostar → equilibrium is needed

1- gravity pulls gas and dust inward

the core



2- $T(\text{core}) \nearrow$

3- $\rho(\text{core}) \nearrow$

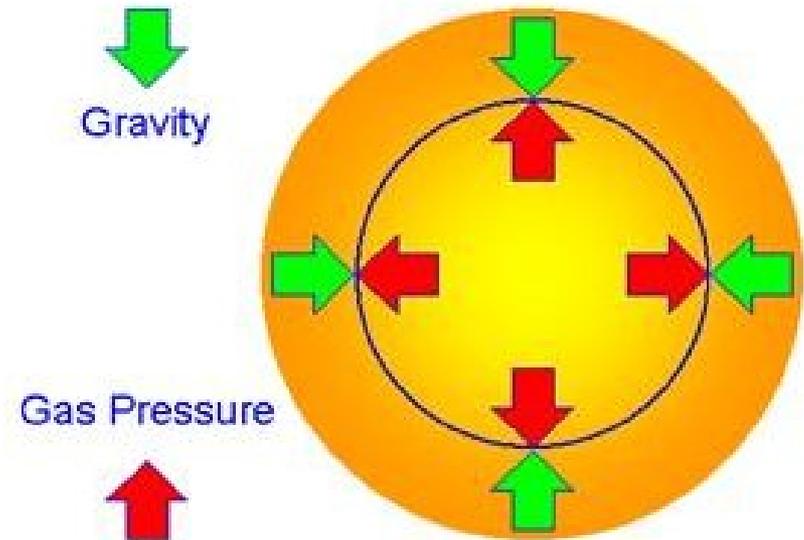


4- gas pressure \nearrow → resists the collapse of the nebulae

5- when gas pressure $(T, \rho) =$ gravity



equilibrium → accretion stops



Option 1: if T not high enough → the protostar ends up a brown dwarf

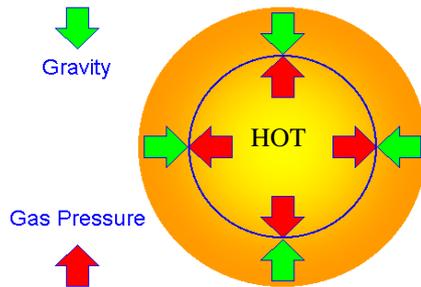
Hydrostatic stellar nucleosynthesis :

Hydrogen burning

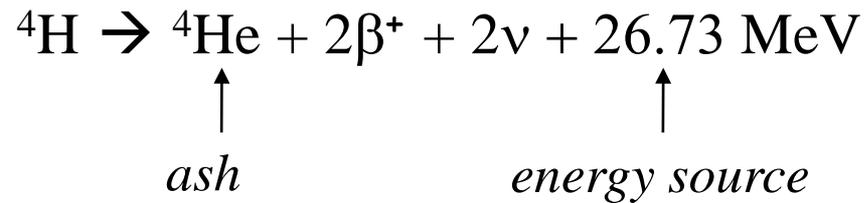
Option 2: if T high enough \longrightarrow “nuclear burning” takes place

when $T \sim 10 - 15 \times 10^6 \text{ K}$ and $\rho \sim 10^2 \text{ gcm}^{-3}$

Hydrostatic Equilibrium



HYDROGEN BURNING (1st equilibrium) \rightarrow a star is born



$$Q = 4(\Delta m)1\text{H} - (\Delta m)4\text{He} \approx 26.73\text{MeV}$$

gravitational collapse is halted \Rightarrow star undergoes phase of *hydrostatic equilibrium*



MAIN SEQUENCE STARS

Minimum mass required for hydrogen burning to ignite $M \sim 0.1 M_{\odot}$

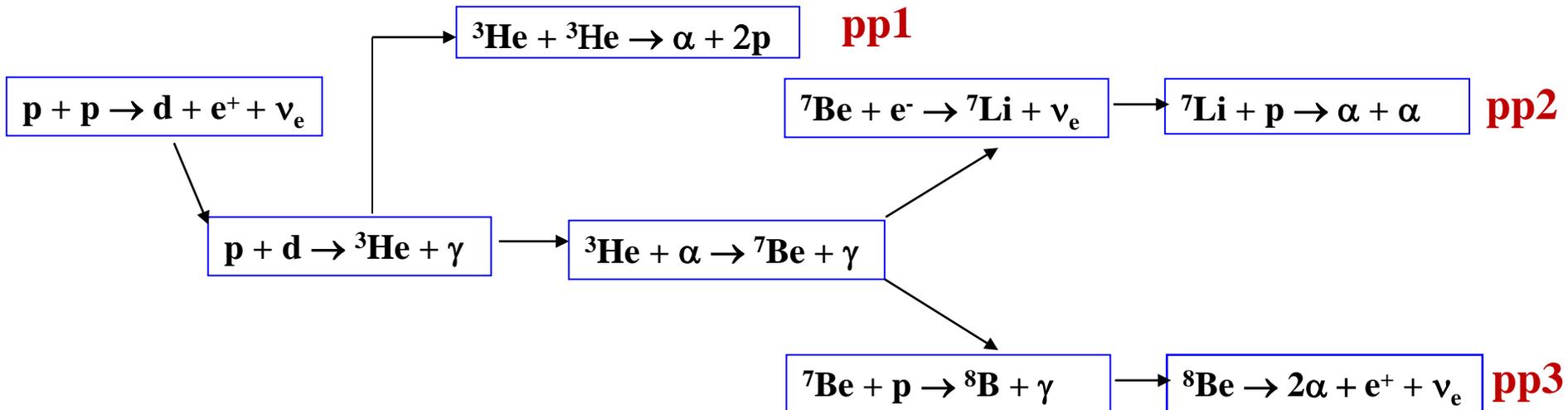
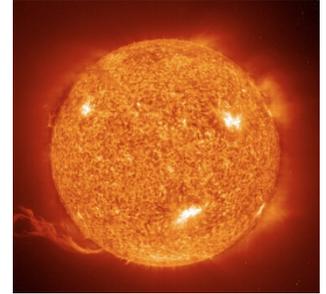
Hydrostatic stellar nucleosynthesis :

Hydrogen burning

- Not simultaneous reaction of 4p (too small probability), but proton-proton chains and CNO cycle

proton-proton (pp) chain

Sun ($T=15.6$ MK),
Core of main-sequence stars ($T=8-55$ MK)
Burning shell of AGB stars ($T=45-140$ MK)



- Neutrinos provide direct evidence that nuclear reaction occurs

pp1 chain

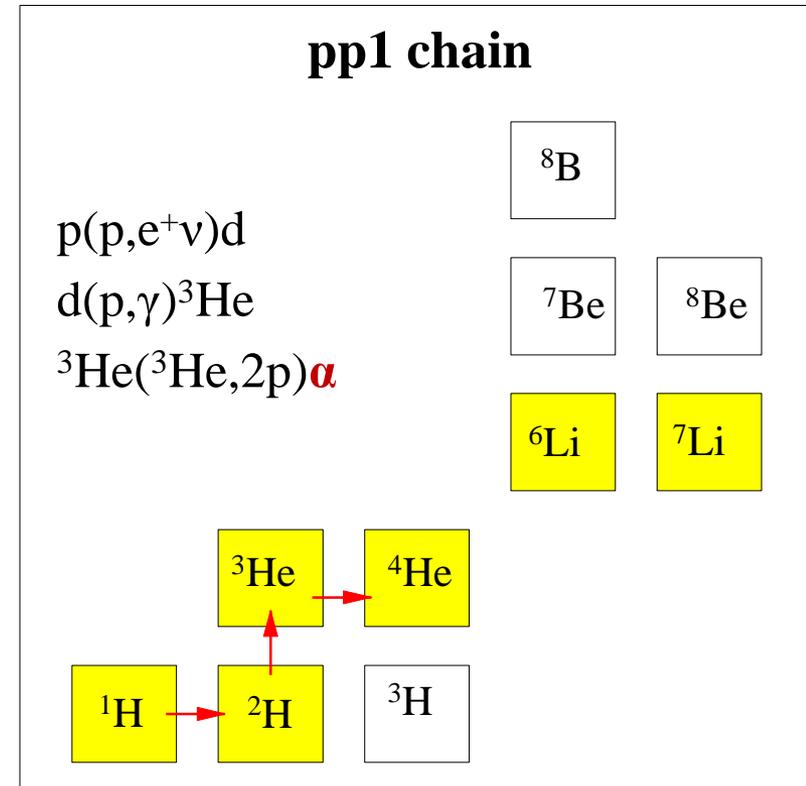
➤ Succession of 3 reactions that produce almost **90%** of Sun's energy

➤ **p+p** ($Q=1.44$ MeV) is a weak interaction
→ the cross-section σ is ~ 20 order of magnitudes smaller than σ associated with nuclear interaction
⇒ **the slowest reaction**
→ it can not be measured
→ theoretically determined

➤ All subsequent reactions involve electromagnetic and nuclear forces
⇒ **much faster**

➤ ${}^2\text{H} + \text{p} \rightarrow {}^3\text{He} + \gamma$ ($Q=5.49$ MeV)
has been measured in LUNA laboratory (lecture IV)

➤ First p+p neutrino detection (BOREXINO) in 2014



pp1 chain

➤ ${}^3\text{He} + \text{p} \rightarrow {}^4\text{Li} \rightarrow {}^3\text{He} + \text{p}$ ($\tau = 10^{-22}$ s)

➤ ${}^3\text{He} + \text{d}$ negligible given the low N_{d}

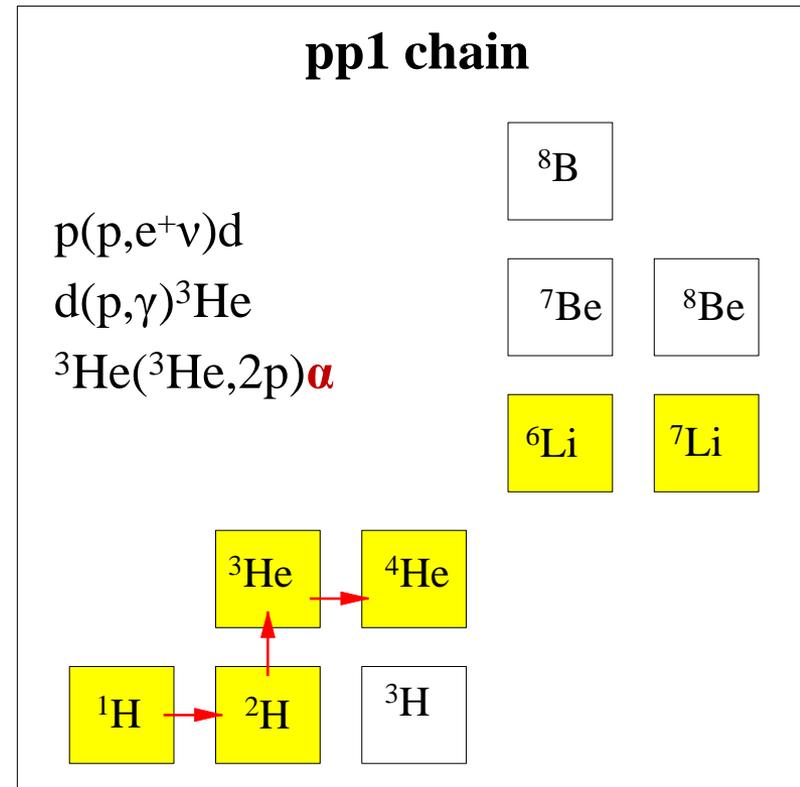
➤ ${}^3\text{He} + {}^3\text{He} \rightarrow \alpha + 2\text{p}$ ($Q=12.86$ MeV)

→ third reaction of the pp1 chain

→ Has been measured in LUNA (lecture IV)

if $N_{\alpha} \gg N_{3\text{He}}$, ${}^3\text{He} + \alpha \rightarrow {}^7\text{Be} + \gamma$

⇒ pp2 and pp3 chains



Hydrostatic stellar nucleosynthesis :

Hydrogen burning

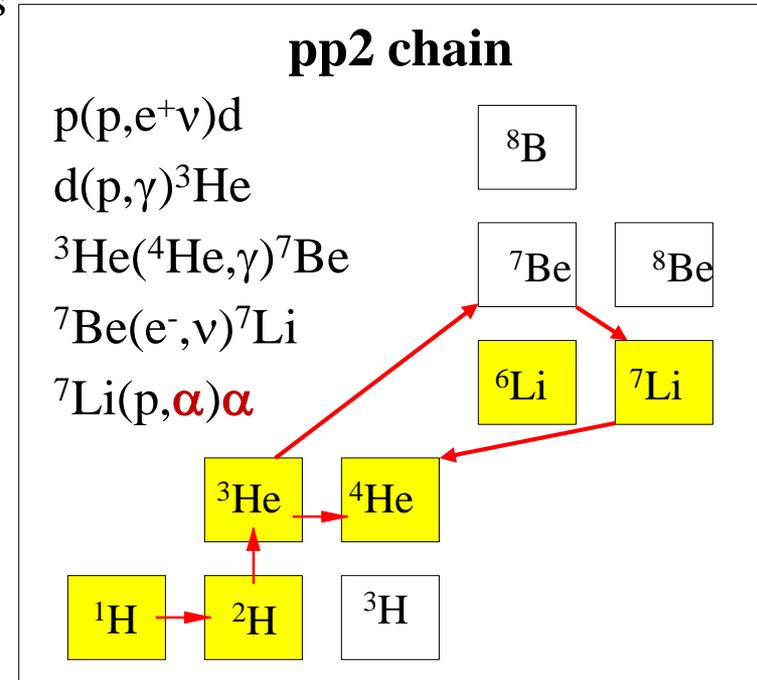
pp2 chain

- ${}^7\text{Be}$ decays by electron capture and its lifetime depends on its charge state
 \Rightarrow if ${}^7\text{Be}$ has all of the orbital electrons removed, the decay rate $\rightarrow 0$
- In center of stars, ${}^7\text{Be}$ almost completely ionized but immersed in a sea of free electrons \Rightarrow the e^- capture rate ($\lambda=1/\tau$) depends on the electron density n_e and T . The stellar lifetime is given by:

$$\tau_s({}^7\text{Be}) = 7.06 \times 10^8 \frac{T_6^{1/2}}{\rho(1 + X_H)}$$

Ex: In the sun's core $T_6=15$, $\rho=100 \text{ g cm}^{-3}$, $X_H=0.5 \rightarrow \tau_s({}^7\text{Be})= 140 \text{ d}$

- But the probability that ${}^7\text{Be}$ atoms are only partially ionized is non-zero
 $\Rightarrow \lambda_{\text{tot}} = \lambda_c + \lambda_K$ (λ_c : decay rate due to continuum free electrons, λ_K : K-capture rate)
 $\Rightarrow \tau_{\text{sun}}({}^7\text{Be}) = \mathbf{120 \text{ d} = 0.33 \text{ y}}$ [$\tau_e({}^7\text{Be}) = 77 \text{ d}$ on earth]

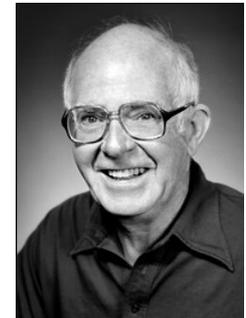
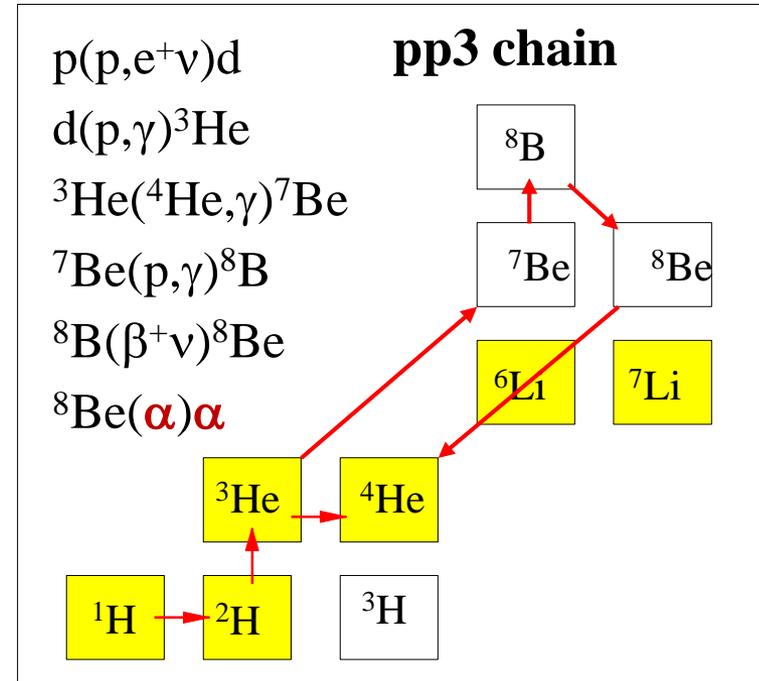


Hydrostatic stellar nucleosynthesis :

Hydrogen burning

pp3 chain

- The enhanced $\tau_s(^7\text{Be})$ in stars increases the probability of its destruction via $^7\text{Be}(p, \gamma)^8\text{Be}$ when T is increasing
- For sun's conditions: T=15MK
 $\tau_s(^7\text{Be}) = 0.33 \text{ y} \ll \tau_H(^7\text{Be}) = 150 \text{ y}$
→ 99.89% of the time, ^7Be proceeds through e- capture & only 0.11% through proton capture
- The reaction $^7\text{Be}(p, \gamma)^8\text{Be}$ is faster for **T > 25 MK**
⇒ **pp3 chain**
- ^8B neutrinos discovered at Homestake [0.02%] in 1968
(solar ν problem)
- Super-Kamiokande/SNO experiments (**neutrino oscillations**)

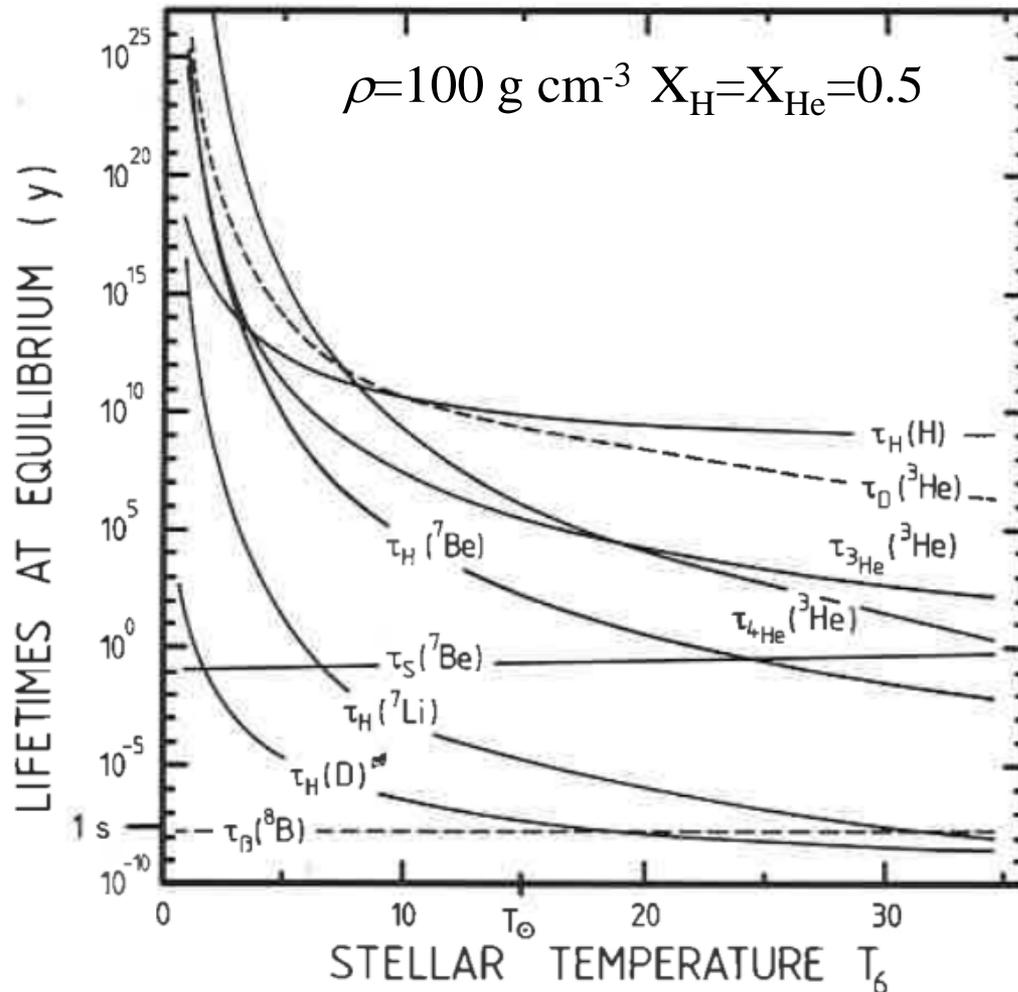


Ray Davis (1914-2006)
Nobel Prize 2002

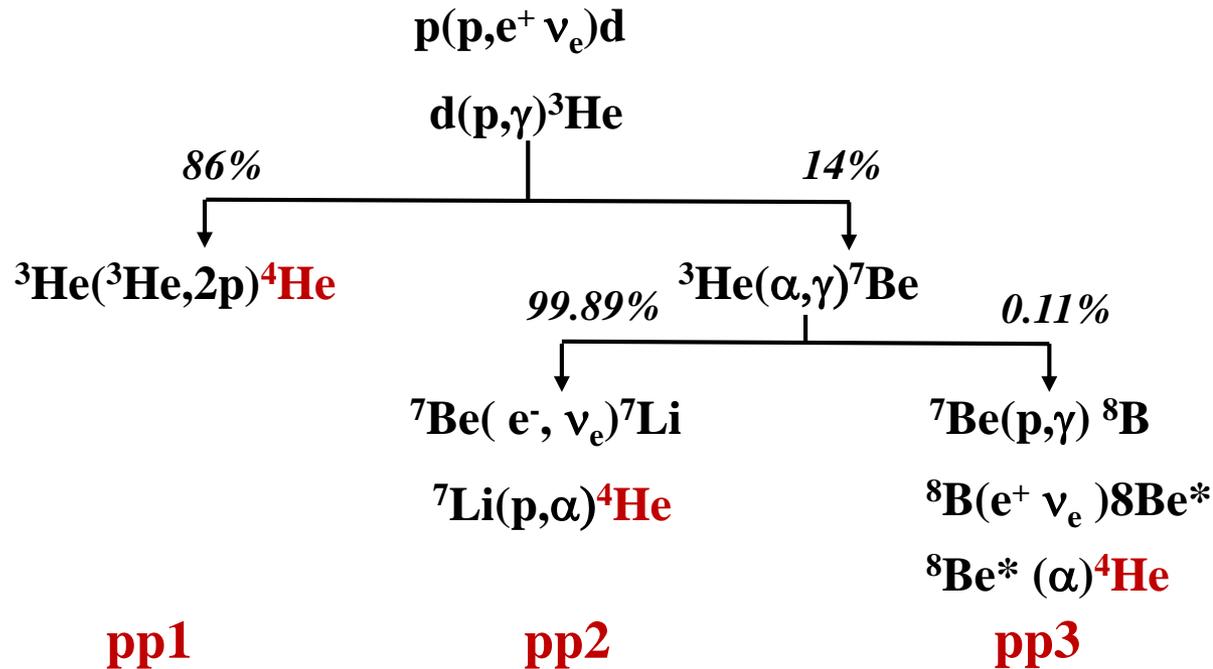
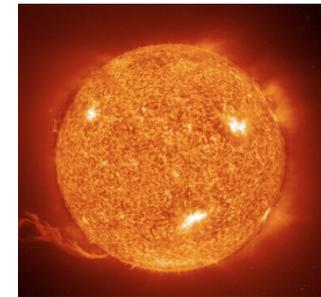
Hydrostatic stellar nucleosynthesis :

Hydrogen burning

pp chains: Capture & decay lifetimes as function of stellar temperature
(see later in lecture III)



The reactions of p-p chain in the sun



pp1-chain:

$$\bar{E}_\nu(p + p) = 0.26 \text{ MeV}$$

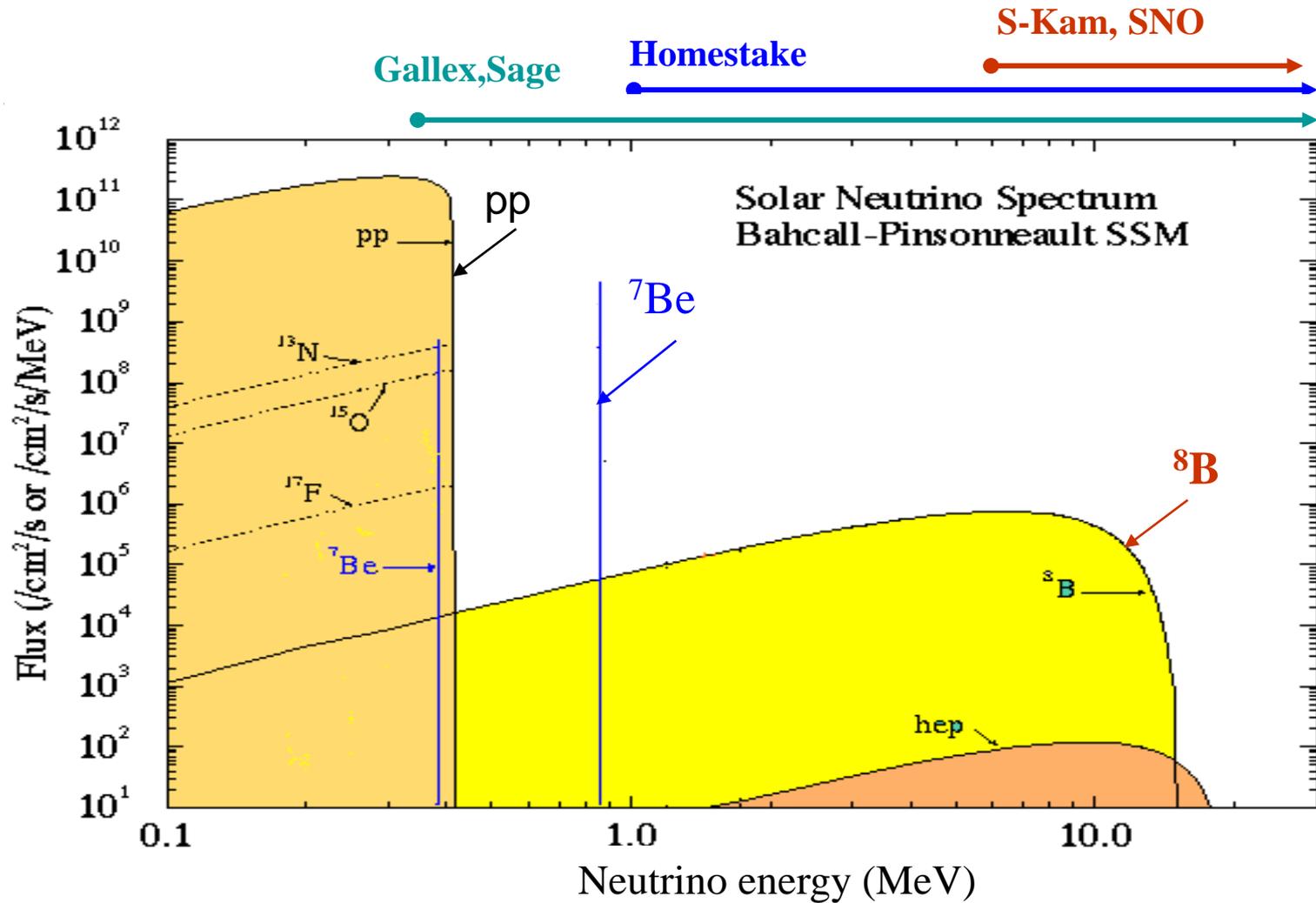
$$Q_{\text{eff}} = Q - 2\bar{E}_\nu(p + p) = 26.20 \text{ MeV}$$

pp2-chain: $\bar{E}_\nu({}^7\text{Be}) = 0.81 \text{ MeV}, Q_{\text{eff}} = Q - 2\bar{E}_\nu(p + p) - \bar{E}_\nu({}^7\text{Be}) = 25.66 \text{ MeV}$

pp3-chain: $\bar{E}_\nu({}^8\text{B}) = 7.30 \text{ MeV}, Q_{\text{eff}} = Q - 2\bar{E}_\nu(p + p) - \bar{E}_\nu({}^8\text{B}) = 19.17 \text{ MeV}$

The effective energy given to the star is $< Q = 26.73 \text{ MeV}$ because of the **escape of the neutrinos**

The solar neutrino problem



- SSM: “Standard” Solar Model

The solar neutrino problem: detection of solar neutrinos (1)

➤ The pioneering experiment of Ray Davis (1968-2001) & J.N. Bahcall

680 tons of perchloroethylene (C_2Cl_4) in the Homestake gold mine (1.5 km deep)



Production of ${}^{37}Ar$: ~0.5 atom per day

Radiochemical separation: extraction of the ${}^{37}Ar$ nuclei every 100 days, counting (EC) in a gas detector

Result: 2.56 ± 0.16 (stat) ± 0.16 (syst) SNU \Rightarrow 30% of the expected signal

Solar model (Bahcall 2004): 8.5 ± 0.18 SNU

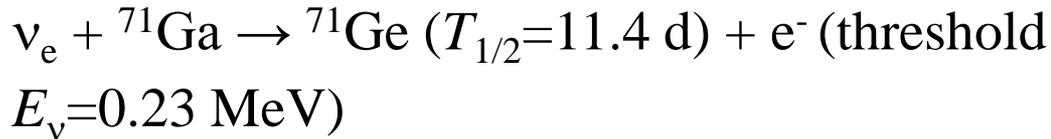


1 SNU (Solar Neutrino Unit) = 10^{-36} capture per second and target atom

The solar neutrino problem: detection of solar neutrinos (2)

- Radiochemical experiments with gallium:

SAGE and GALLEX Reaction:



→ Result: 40% of the expected signal

- Real-time detection of (mostly) **e^- neutrinos**:

Kamiokande (700 t of water, 1983 - 1996),

Super-Kamiokande (50 kt, 1996 → M.

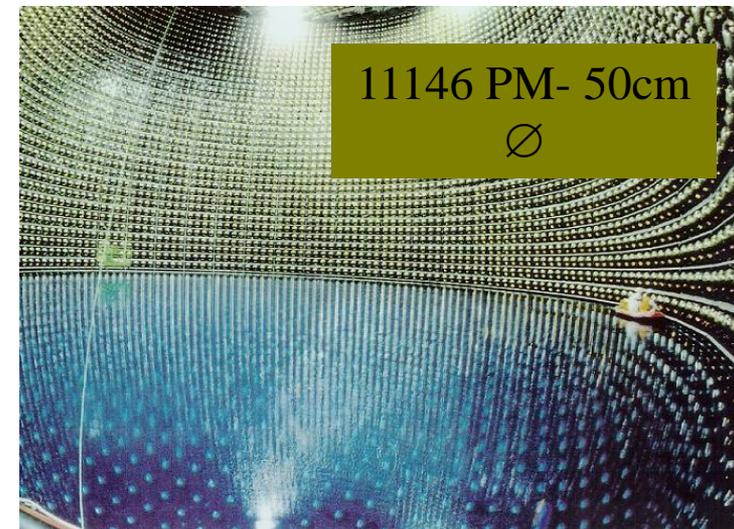
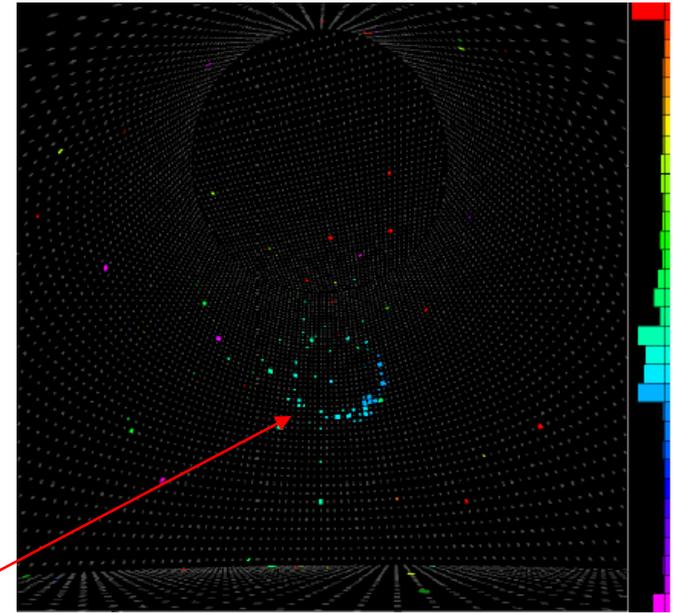
Koshiya Nobel price in 2002) $\nu_e + e^- \rightarrow e^- + \nu_e$

⇒ **Cherenkov light**

Result: $2.44 \pm 0.102(\text{stat}) \pm 0.08 (\text{syst}) \times 10^6$
 $\text{cm}^{-2}\text{s}^{-1} \Rightarrow 42\%$ of the expected signal

Solar model (Bahcall 2004):

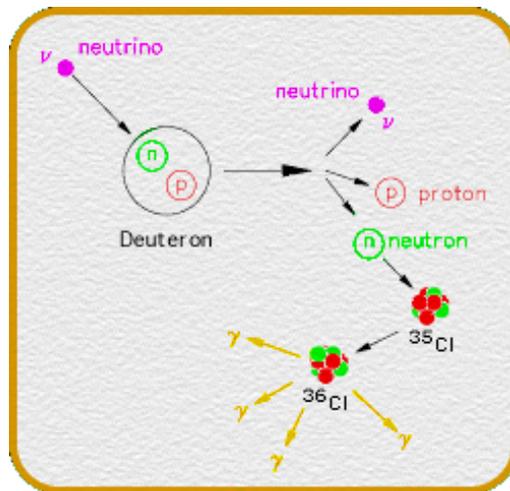
$$5.8 \pm 1.3 \times 10^6 \text{ cm}^{-2}\text{s}^{-1}$$



The solar neutrino problem:

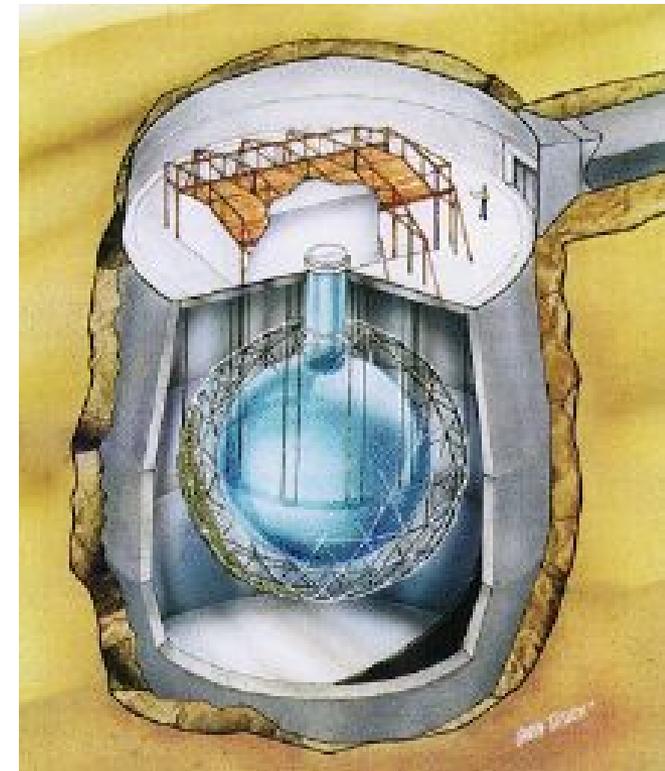
The solution

- Origin of the deficit: problem with the **nuclear data: ${}^7\text{Be}(p,\gamma){}^8\text{B}$?**, or new physics of the neutrino = **oscillation $\nu_e \rightarrow \nu_\mu, \nu_\tau$?**
- Sudbury Neutrino Observatory (SNO): 1100 tons of D_2O (99.9%) **sensitive to the 3 neutrino flavours** $\nu_x + d \rightarrow p + n + \nu_x$ (neutral current)



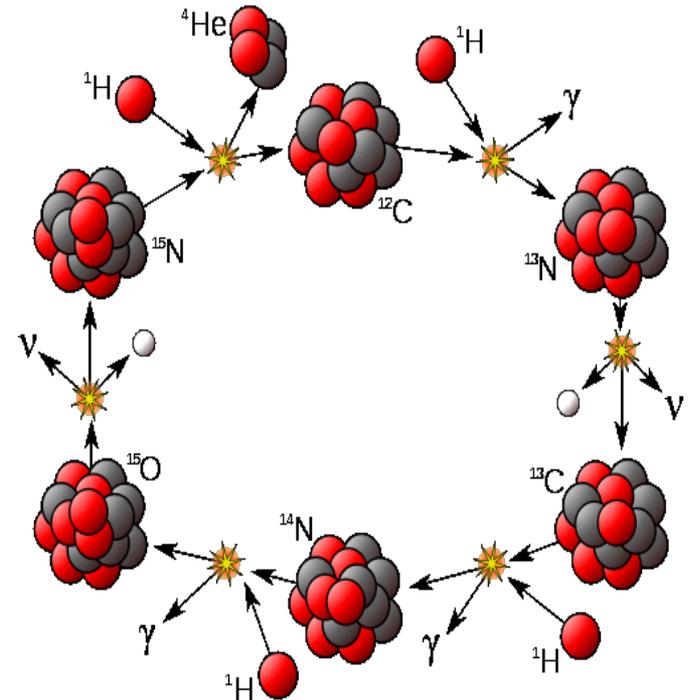
$$\phi_{\text{CN}} = 5.21 \pm 0.27 \text{ (stat)} \pm 0.389 \text{ (syst)} \times 10^6 \text{ cm}^{-2}\text{s}^{-1}$$

in **agreement** with $\phi_{\text{SSM}} = 5.05^{+1.01}_{-0.81} \times 10^6 \text{ cm}^{-2}\text{s}^{-1}$



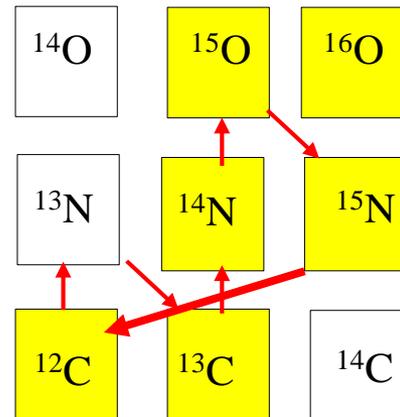
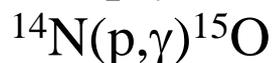
The CNO cycle

- In population I stars ([second-or third-generation stars](#)), the elements C, N and O can serve as **catalysts** of the transformation $4p \rightarrow {}^4\text{He} + 2e^+ + 2\nu$ ($Q = 26.73 \text{ MeV}$)



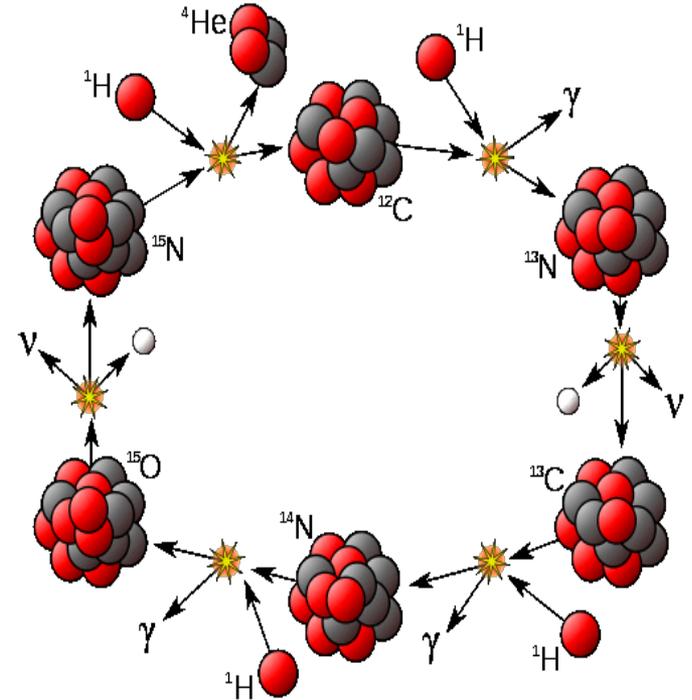
- If the cycle starts with ${}^{12}\text{C}$, it ends with ${}^{12}\text{C} \Rightarrow {}^{12}\text{C}$ is a catalyst & only hydrogen is consumed

CNO-1 cycle



The CNO cycle

- The rate of energy production in the CNO cycle $\epsilon \propto Q_{\text{CNO}} / \tau_{\text{CNO}}$ as well as the **cycle duration** is governed by the **slowest** reaction $^{14}\text{N}(p,\gamma)^{15}\text{O}$: $\tau_{\text{CNO}} = \tau_p(^{12}\text{C}) + \tau_p(^{13}\text{C}) + \tau_p(^{14}\text{N}) + \tau_p(^{15}\text{N}) \cong \tau_p(^{14}\text{N})$
 $(\tau_{\beta^+}(^{13}\text{N}) = 14 \text{ min} \ \& \ \tau_{\beta^+}(^{15}\text{O}) = 2.9 \text{ min})$
 (being negligible)



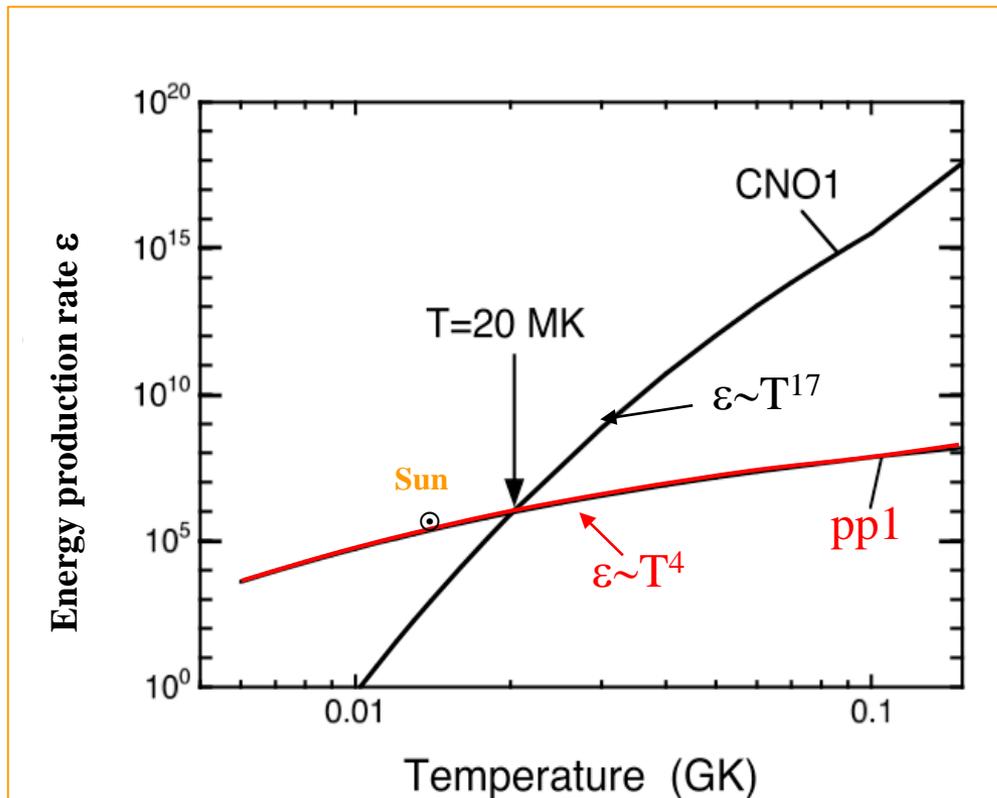
Ex: for $\rho=100 \text{ g/cm}^3$, $X_{\text{H}}=0.5$ & $T=60 \text{ MK}$

$\tau_p(^{12}\text{C})=6.1 \times 10^9 \text{ y}$; $\tau_p(^{13}\text{C})=1.1 \times 10^9 \text{ y}$; $\tau_p(^{14}\text{N})=2.1 \times 10^{12} \text{ y}$; $\tau_p(^{15}\text{N})=1.0 \times 10^8 \text{ y}$

- $^{14}\text{N}(p,\gamma)^{15}\text{O}$ has been measured by LUNA (lecture IV)
- CNO cycle (in AGB stars) = main source of ^{13}C and ^{14}N in the Universe

The CNO cycle

CNO : steeper temperature–dependence than pp chain (see later in lecture III)
⇒ pp chain dominates in low–mass main–sequence stars ($\sim M_{\odot}$) and CNO cycle dominates in higher–mass stars (few M_{\odot})

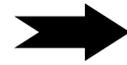


- $T > 20$ MK: CNO1 faster than pp1
 - $M \geq 1.5 M_{\odot} \Rightarrow T_6 > 30$
↓
CNO cycle
(also depends on CNO abundance)

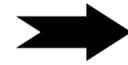
Hydrostatic stellar nucleosynthesis :

Helium burning

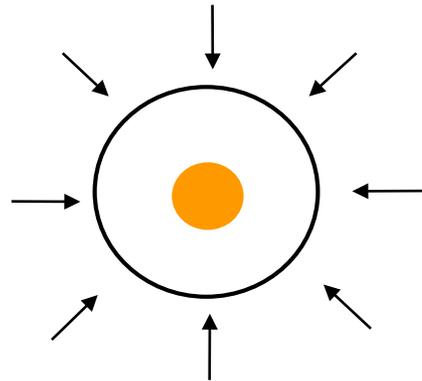
H exhausted in core
isothermal He core
contraction sets in



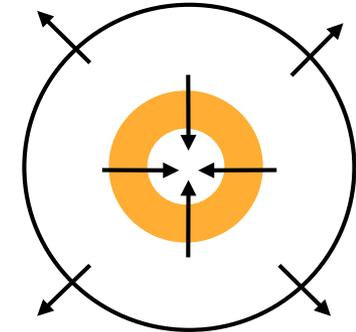
temperature
increases



H-burning shell



$$R \sim 10-100 R_{\odot} \Rightarrow T_s \sim 3-4 \times 10^3 \text{ K}$$



RED GIANT STARS

First step in old age

contracting core
expanding envelope

when $T \sim 10^8 \text{ K}$ and $\rho \sim 10^3 \text{ gcm}^{-3}$ (minimum mass $\sim 0.5 M_{\odot}$)



HELIUM BURNING

(2nd equilibrium)

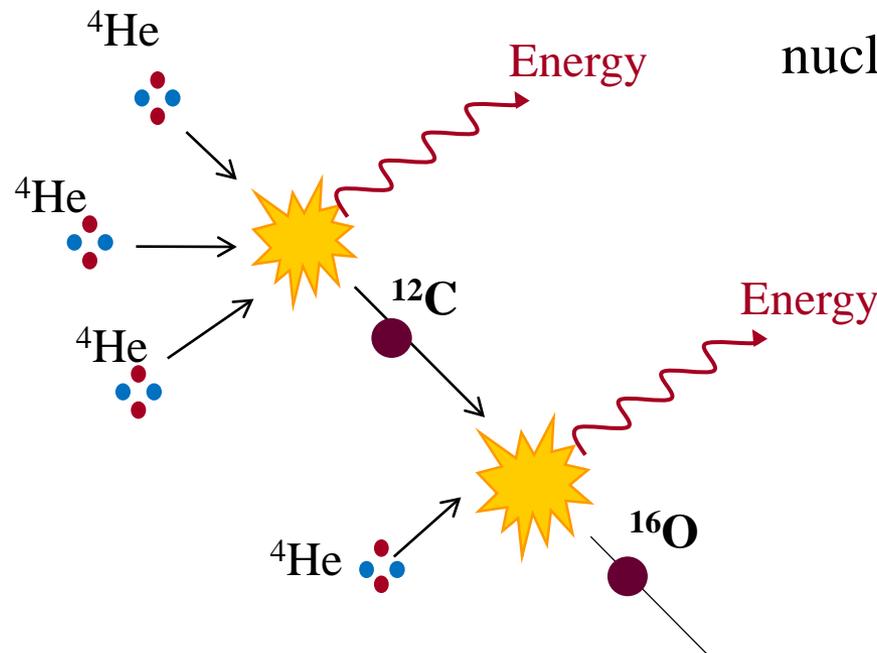
Hydrostatic stellar nucleosynthesis :

Helium burning

- Synthesis of ^{12}C and ^{16}O from ^4He in stars of more than $\sim 0.5 M_{\odot}$

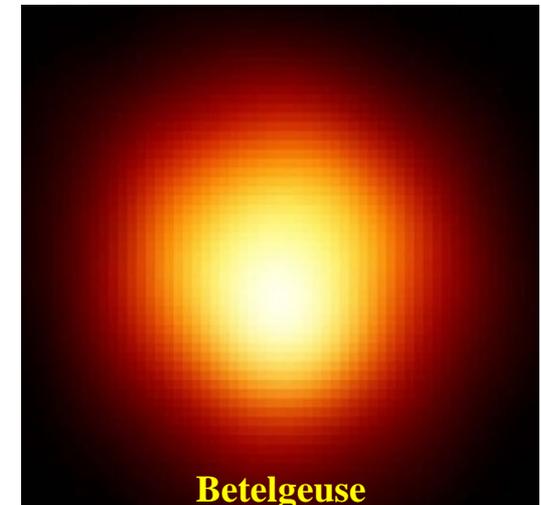
$$T_{\text{C}} = (1 - 4) \times 10^8 \text{ K}$$

Main reactions:



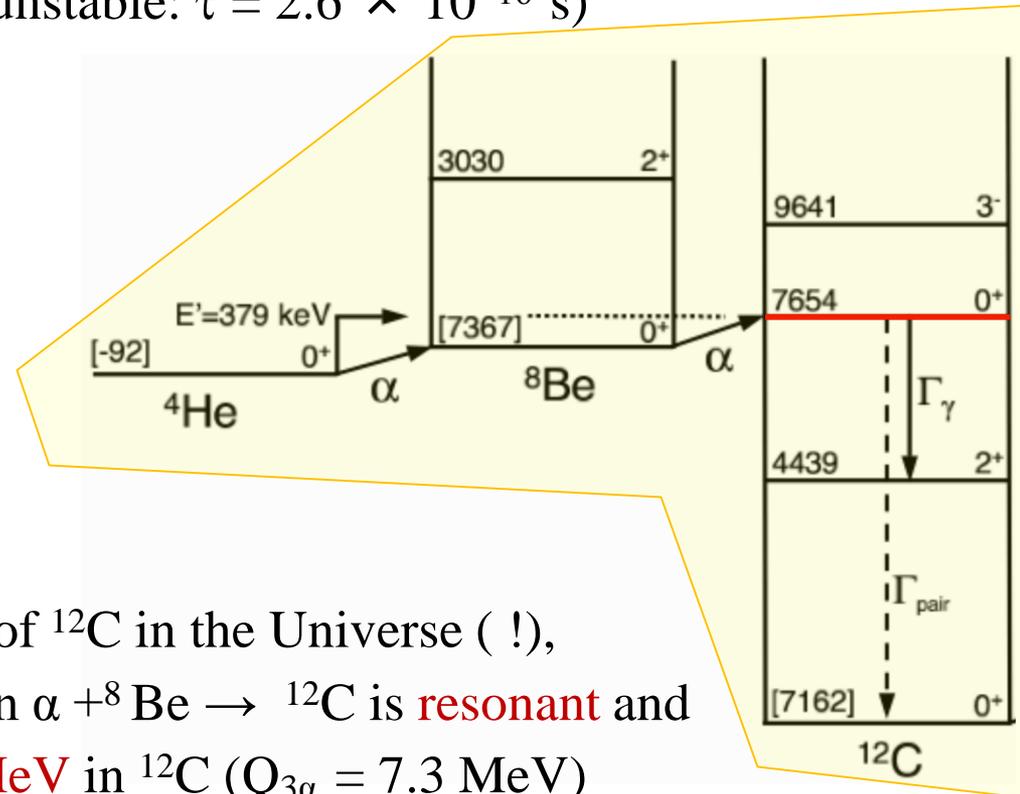
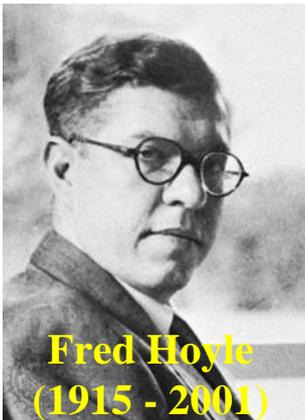
nuclear burning ashes

energy source

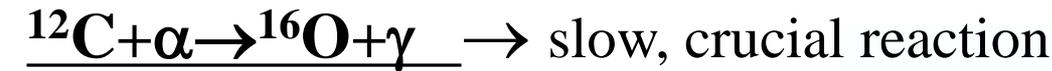


The triple α -process:

- How are synthesized elements heavier than ${}^4\text{He}$, given that there are **no stable isotopes** for mass $A=5$ ($p+\alpha$) and $A=8$ ($\alpha+\alpha$)?
- Fusion of 3α in ${}^{12}\text{C}$ in **2 steps**:
 $\alpha + \alpha \rightarrow {}^8\text{Be}$ $Q=-92$ keV (${}^8\text{Be}$ is unstable: $\tau = 2.6 \times 10^{-16}$ s)
then $\alpha + {}^8\text{Be} \rightarrow {}^{12}\text{C}^*$ $Q=-288$ keV



- In view of the significant abundance of ${}^{12}\text{C}$ in the Universe (!), Hoyle (1953) **predicted** that the reaction $\alpha + {}^8\text{Be} \rightarrow {}^{12}\text{C}$ is **resonant** and the **existence of a $J^\pi = 0^+$ state at 7.7 MeV** in ${}^{12}\text{C}$ ($Q_{3\alpha} = 7.3$ MeV)
- Experimental verification in 1953 and 1957

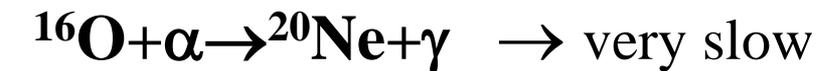


- The rate of this reaction determines the $^{12}\text{C}/^{16}\text{O}$ ratio at the end of the helium burning phase (solar abundance $(\text{C}/\text{O})_{\odot} = 0.5$).

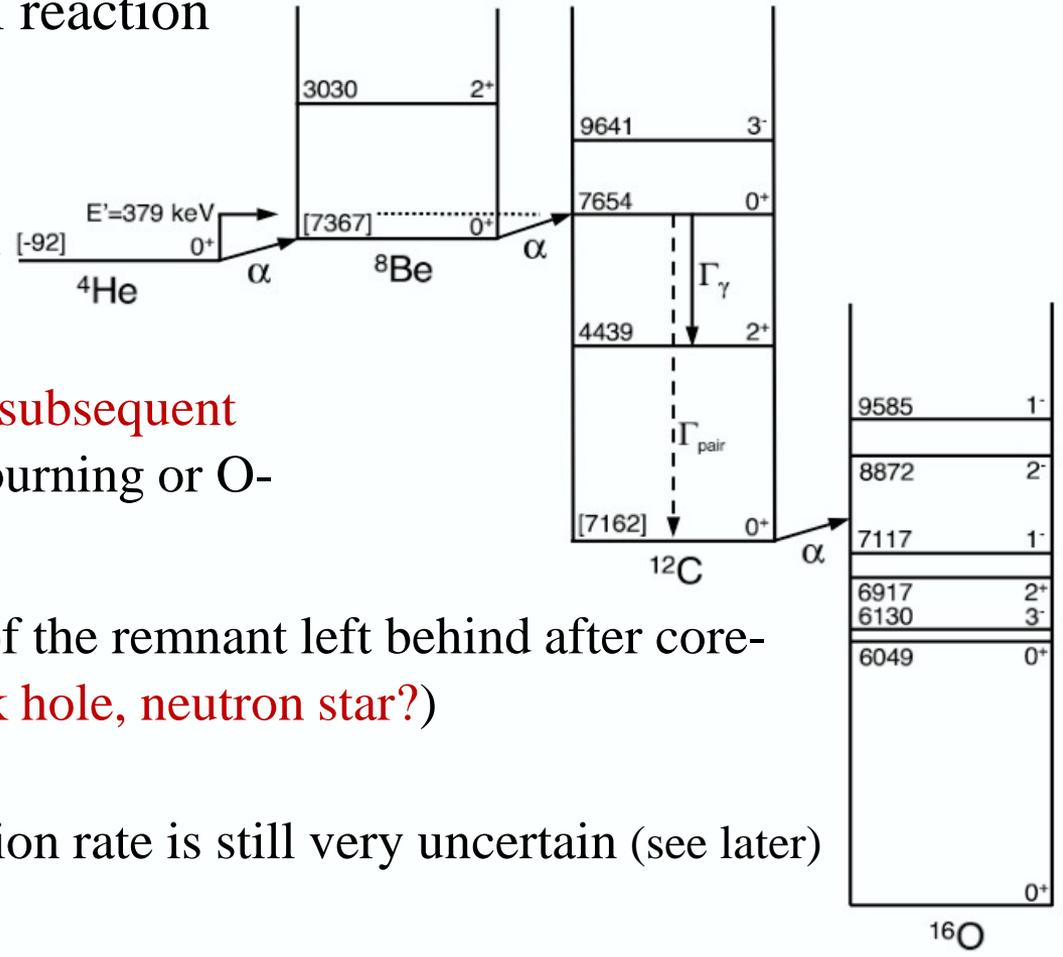
- $^{12}\text{C}/^{16}\text{O}$ ratio has impact on the subsequent burning stages in massive stars (C-burning or O-burning?)

- $^{12}\text{C}/^{16}\text{O}$ ratio influences the nature of the remnant left behind after core-collapse supernova explosion (black hole, neutron star?)

- Various measurements but the reaction rate is still very uncertain (see later)

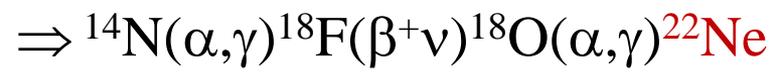


- reaction rate \ll rate of $^{12}\text{C}(\alpha, \gamma)^{16}\text{O}$ for $T < 3 \times 10^8 \text{ K} \Rightarrow$ survival of ^{16}O at the end of helium burning



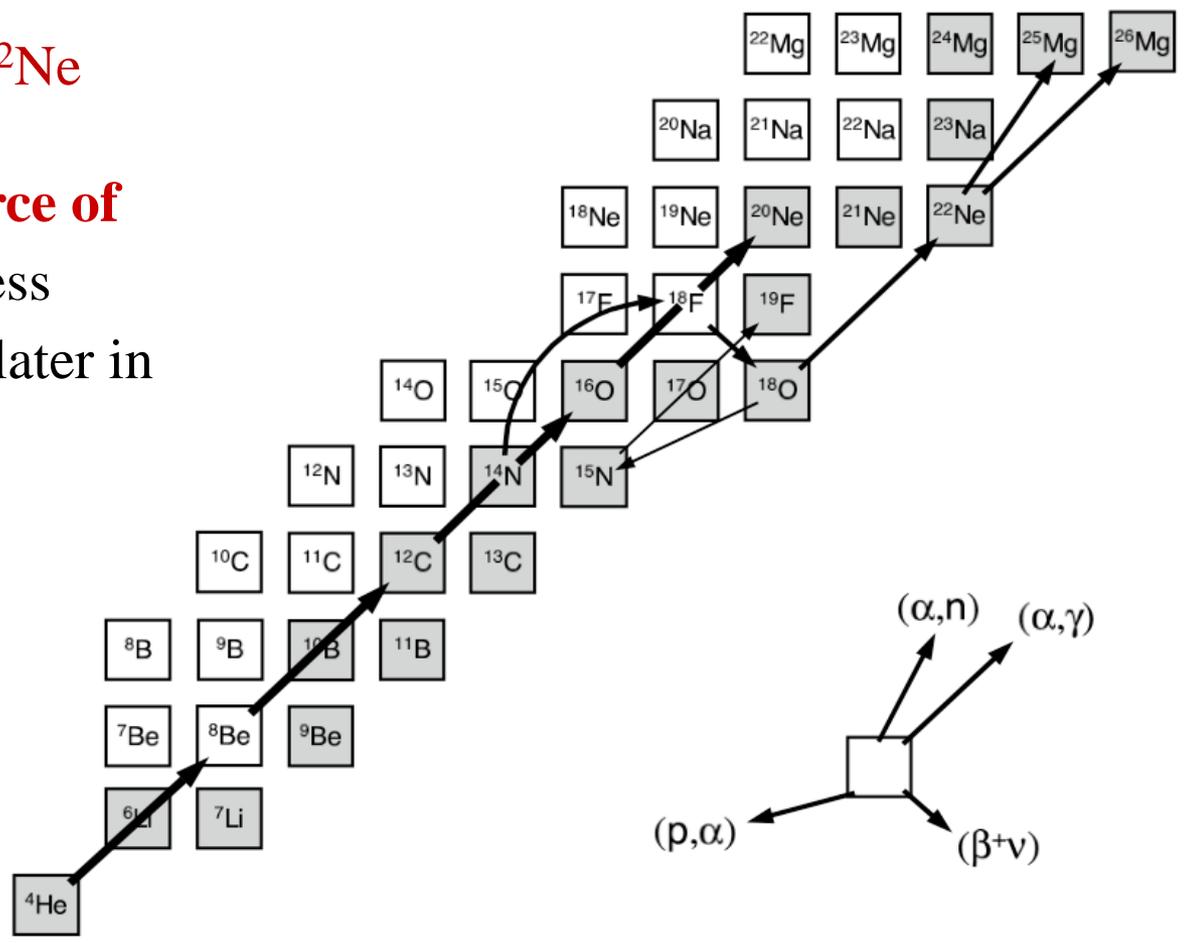
Other reactions:

- ^{14}N is produced by the CNO cycle and accounts for 1-2% of the mass of the fusion core at the end of H burning (pop I stars)



+ $^{22}\text{Ne}(\alpha,n)^{25}\text{Mg}$ = main **source of neutrons** of the weak s-process
(“slow” neutron capture, see later in this lecture)

- Helium burning: main **source** of ^{12}C , ^{16}O , ^{18}O and ^{22}Ne in the Universe



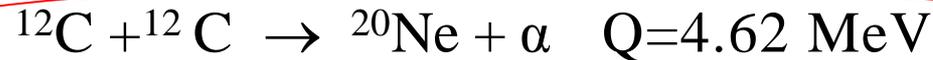
Hydrostatic stellar nucleosynthesis: Advanced burning stages (I)

He exhausted in the core → mainly of ^{12}C & ^{16}O in the core → gravitational contraction
→ increase of T. When $T \sim (5-9) \times 10^8 \text{ K}$ and $\rho > 2 \cdot 10^5 \text{ gcm}^{-3}$ ($M \geq 8 M_{\odot}$)

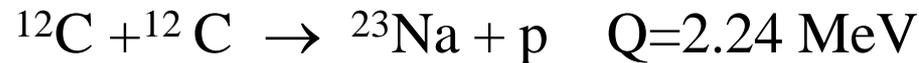


Carbon burning

Major reaction sequences:



dominates by far



+several secondary reactions

Composition at the end of burning:

Mainly ^{20}Ne with some $^{21,22}\text{Ne}$, ^{23}Na , $^{24,25,26}\text{Mg}$, $^{26,27}\text{Al}$

and ^{16}O : still present in quantities comparable with ^{20}Ne (not burning ... yet)

Hydrostatic stellar nucleosynthesis: Advanced burning stages (II)

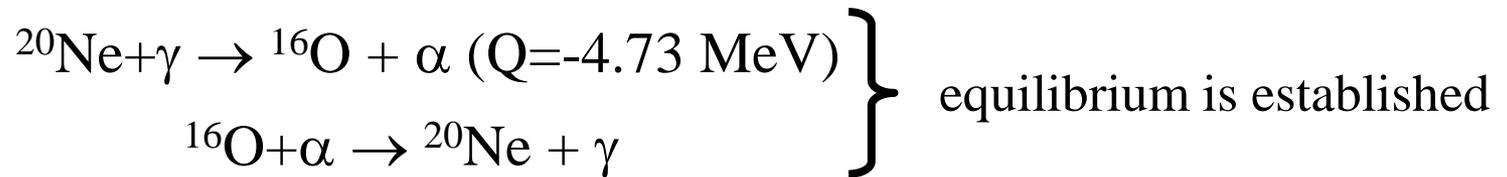
➤ After C burning, when $T \sim (1-2) \times 10^9 \text{ K}$ and $\rho \sim 10^6 \text{ gcm}^{-3}$ ($M \geq 11 M_{\odot}$)

➔ Neon burning

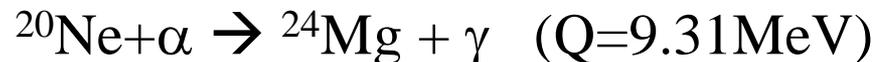
Why would neon burn before oxygen ???

Answer:

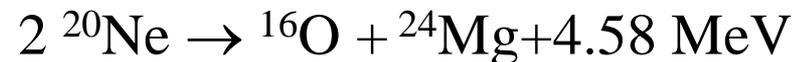
Temperatures are sufficiently high to initiate **photodisintegration** of ^{20}Ne



this is followed by (using the released helium)



so net production of nuclear energy :



➤ Main nuclear ash **^{16}O** with some ^{24}Mg

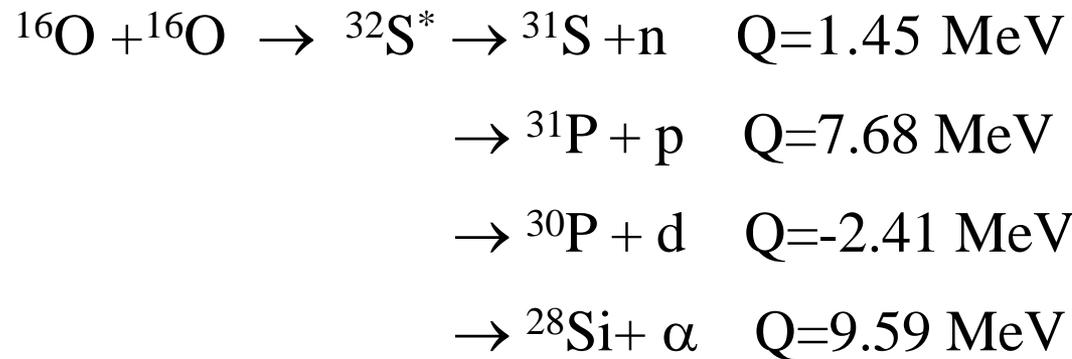
Hydrostatic stellar nucleosynthesis: Advanced burning stages (III)

After Ne burning, the core contracts further until $T \sim (2-3) \times 10^9 \text{ K}$ and $\rho \sim 3 \times 10^6 \text{ g.cm}^{-3}$

($M \geq 11 M_{\odot}$)

➔ Oxygen burning

Major reaction sequences:



+ recapture of n,p,d, α

Main nuclear ashes:

${}^{28}\text{Si}$ and to a somewhat lesser extent ${}^{32}\text{S}$

Hydrostatic stellar nucleosynthesis: Advanced burning stages (IV)

After O burning, the core contracts until $T \sim (3-4) \times 10^9 \text{ K}$ and $\rho \sim 3 \times 10^7 \text{ g.cm}^{-3}$

($M \geq 11 M_{\odot}$)

➔ Silicon burning

- Starts with ^{28}Si photodesintegration: $^{28}\text{Si}(\gamma, \alpha)^{24}\text{Mg}(\gamma, \alpha)^{20}\text{Ne}(\gamma, \alpha) \dots$
- Photodesintegration rearrangements: destruction of less tightly bound species and capture of released p, n, α to synthesize more tightly bound species
- Many reactions achieve **nuclear statistical equilibrium (NSE)**
 - **NSE** is established when both, photodisintegration rates of the type:
 $(Z, N) + \gamma \rightarrow (Z-1, N) + p$; $(Z, N) + \gamma \rightarrow (Z, N-1) + n$; $(Z, N) + \gamma \rightarrow (Z-2, N-2) + \alpha$
and capture reactions of type: $(Z, N) + p \rightarrow (Z+1, N) + \gamma$; $(Z, N) + n \rightarrow (Z, N+1) + \gamma$;
 $(Z, N) + \alpha \rightarrow (Z+2, N+2) + \gamma$ **are fast**
- equilibrium drives nuclei towards **$A=56$** : most stable nuclei (higher binding energy)
- synthesis of nuclei from Si to Zn (“**iron peak**” elements Ti to Zn)
- Main ash is **^{56}Fe** → formation of an **iron core**

After ^{28}Si burning (Melting), the star core is mainly made of iron. The final stage of a massive star has an onion-ring structure like

→ PRE-SUPERNOVA STARS

➤ Further reactions become **endothermic**

➤ Stars can no longer convert mass into energy via nuclear fusion!

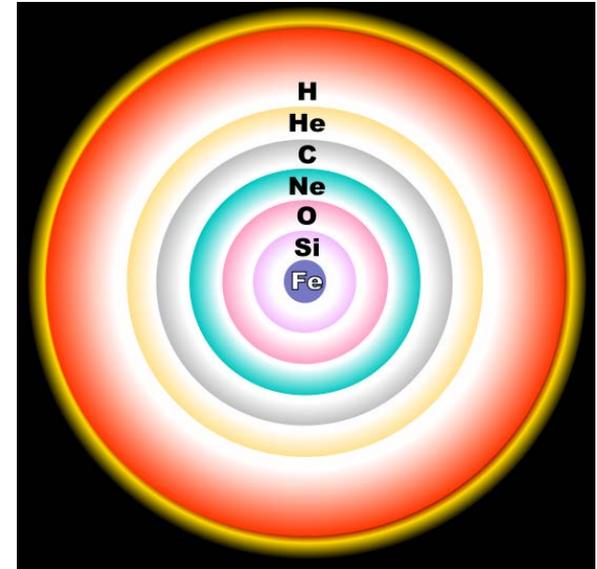
⇒ final gravitational collapse

⇒ **SUPERNOVA EXPLOSION**

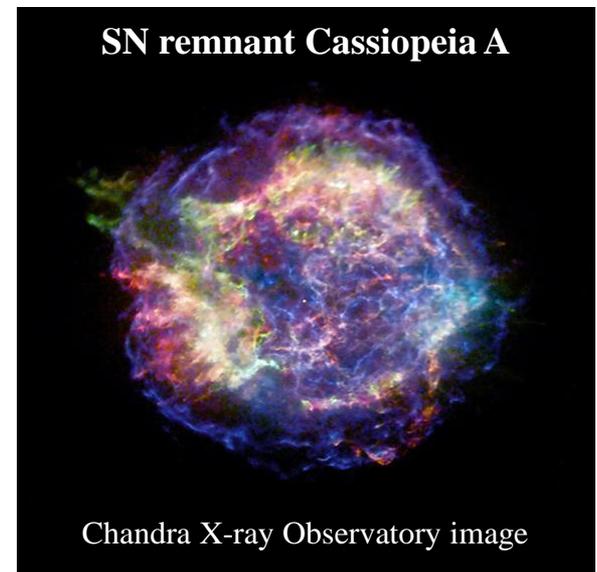
(Core-collapse Supernova or Supernova type II)

⇒ remnant: neutron star or black hole

Note: Burning also takes place in thin regions (**burning shells**) at the interface of different compositional layers



SN remnant Cassiopeia A



Chandra X-ray Observatory image

Stellar mass (M_{\odot})	Stage reached
< 0.08	no thermonuclear fusion
0.1 - 0.5	H burning
0.5 - 8	He burning
8 - 11	C burning
> 11	all stages

Evolution stages of a 25 M_{\odot} star

Stage reached	Timescale	T_{core} (10^9 K)	Density (g cm^{-3})
H burning	7×10^6 y	0.06	5
He burning	5×10^5 y	0.23	7×10^2
C/O burning	600 y / 6 months	0.93 – 2.3	$2 \times 10^5 - 1 \times 10^7$
Si melting	1 d	4.1	3×10^7
Explosive burning	0.1 – 1 s	1.2 - 7	varies

Back to the HR diagram: Main sequence and Red Giant Branch

➤ H-burning core \equiv main sequence
 and H-burning shell \equiv red giant branch (If $M_{\text{star}} \geq 10M_{\odot} \rightarrow$ red supergiant)

➤ H-burning core leads to He core:

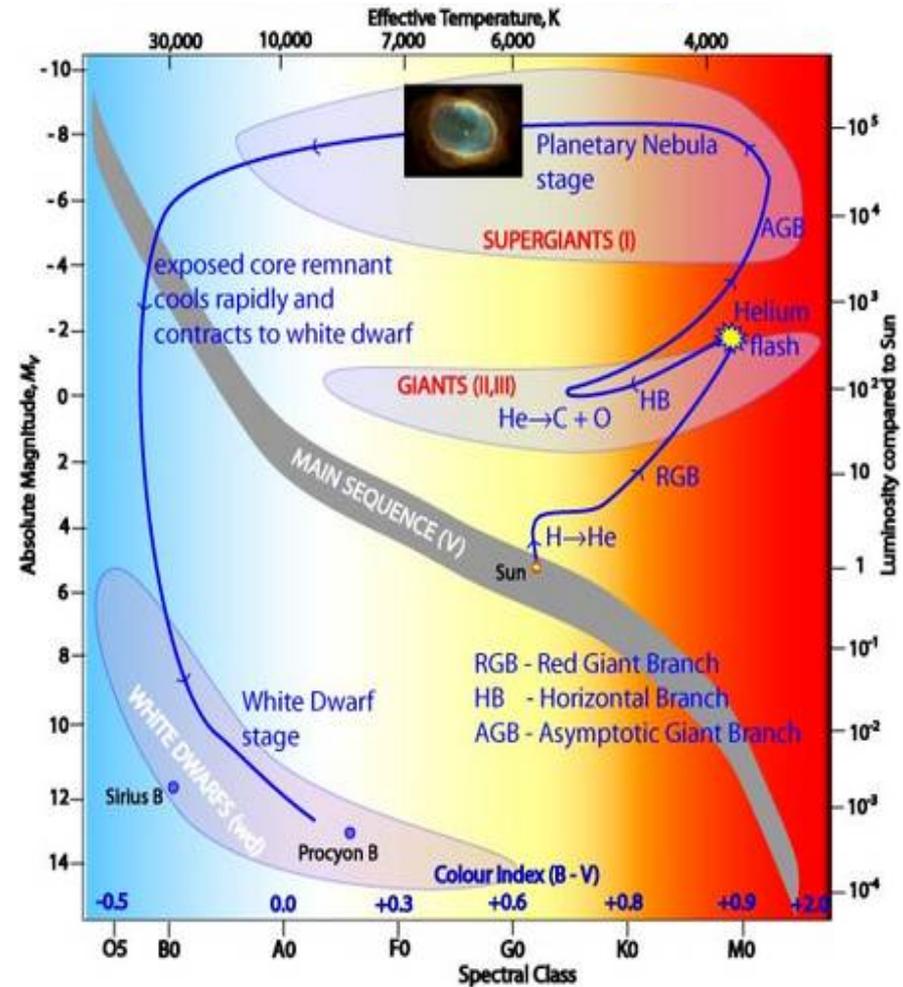
$M_c \nearrow$ core contracts, $T_c \nearrow$, $P \nearrow$

\Rightarrow He inert core surrounded by H-burning shell \Rightarrow expansion and cooling of the envelope

($R_{\odot} \rightarrow 50 R_{\odot} \sim$ Mercury)

➤ subgiant phase (horizontal in HR diagram : $R \nearrow$, $T_{\text{eff}} \searrow$ and $L \approx \text{cst}$)

➤ red giant phase (vertical in HR diagram : $L \nearrow$ at $T_{\text{eff}} \searrow$)

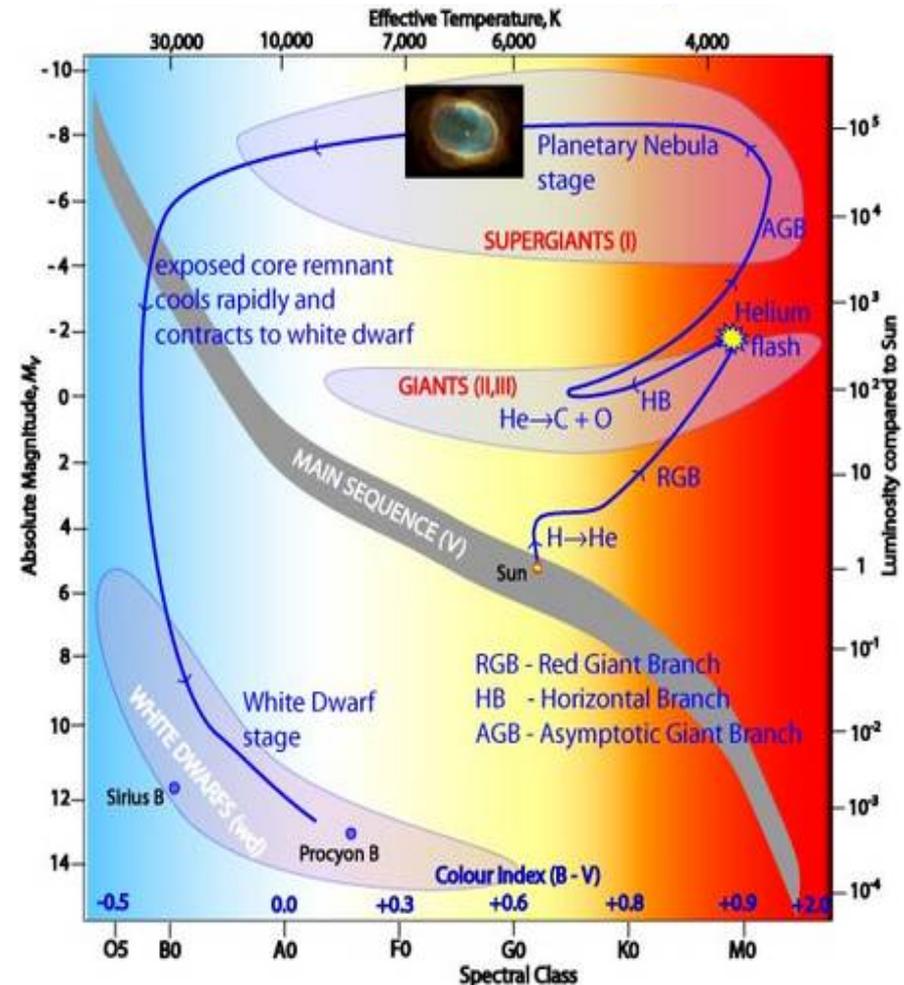


- **Ignition of the He core** at $T_C \sim 10^8$ K \Rightarrow core contraction stops

- In low-mass stars (0.7-2 M_\odot) the electron gas in the centre is partially **degenerate** \Rightarrow **helium flash**

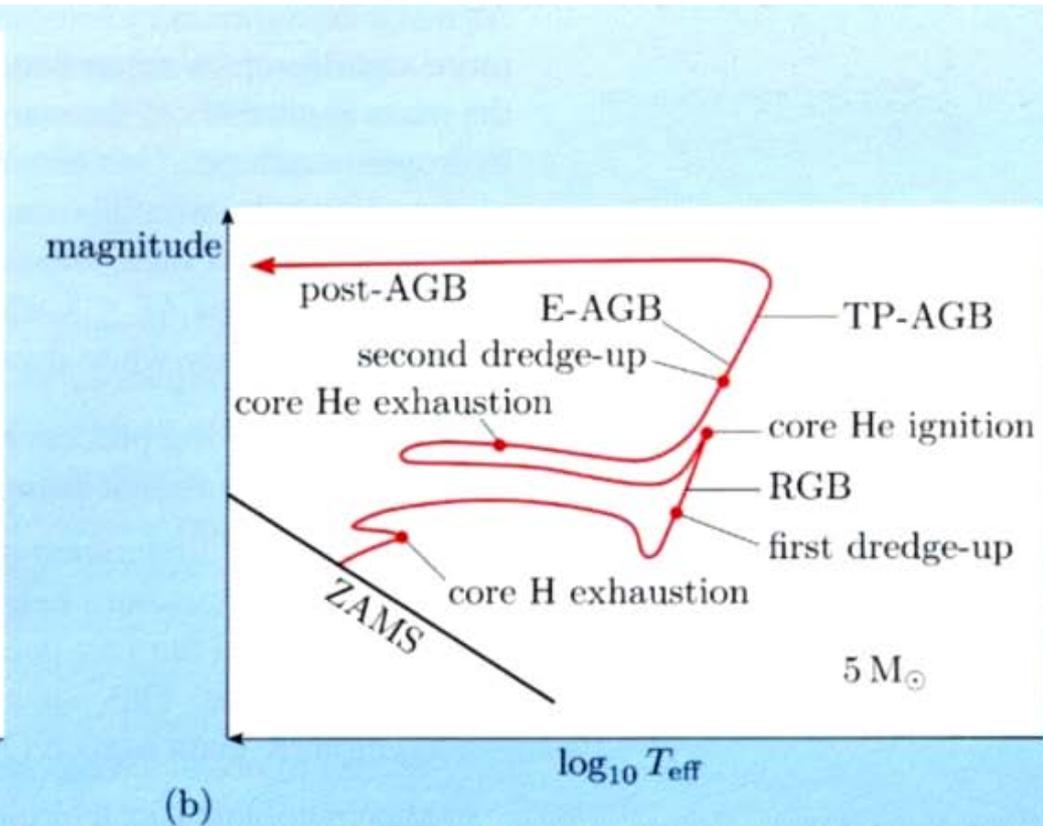
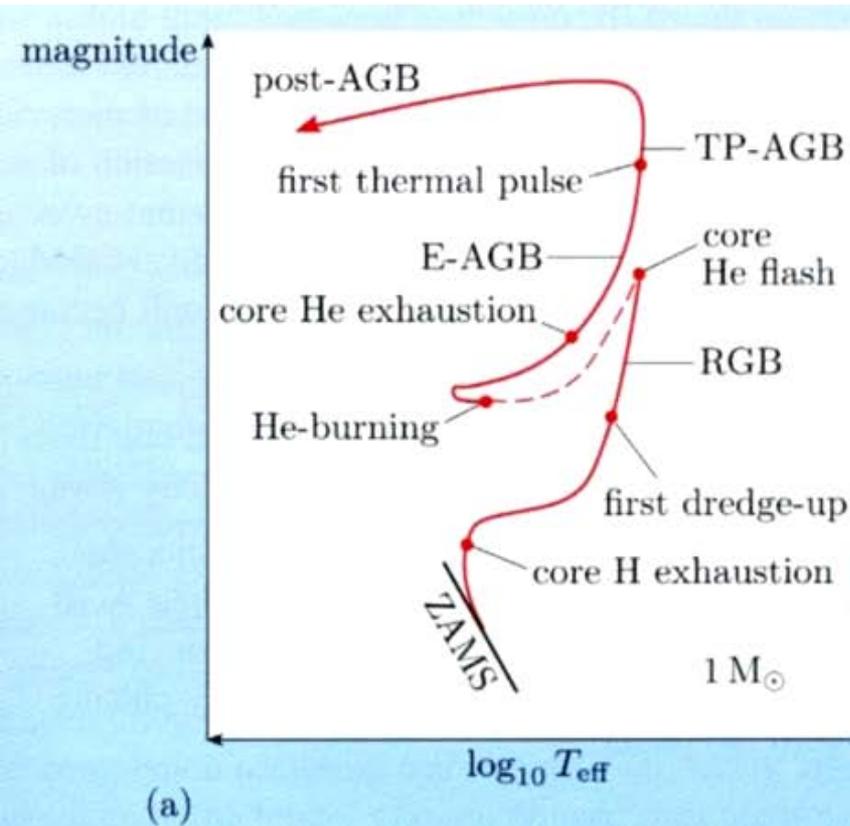
\Rightarrow Release during a few seconds of $10^{10} L_\odot$ in L_{nuc} !
But **invisible from the surface**

\Rightarrow Expansion and cooling of the core result in the **contraction and heating of the envelope**



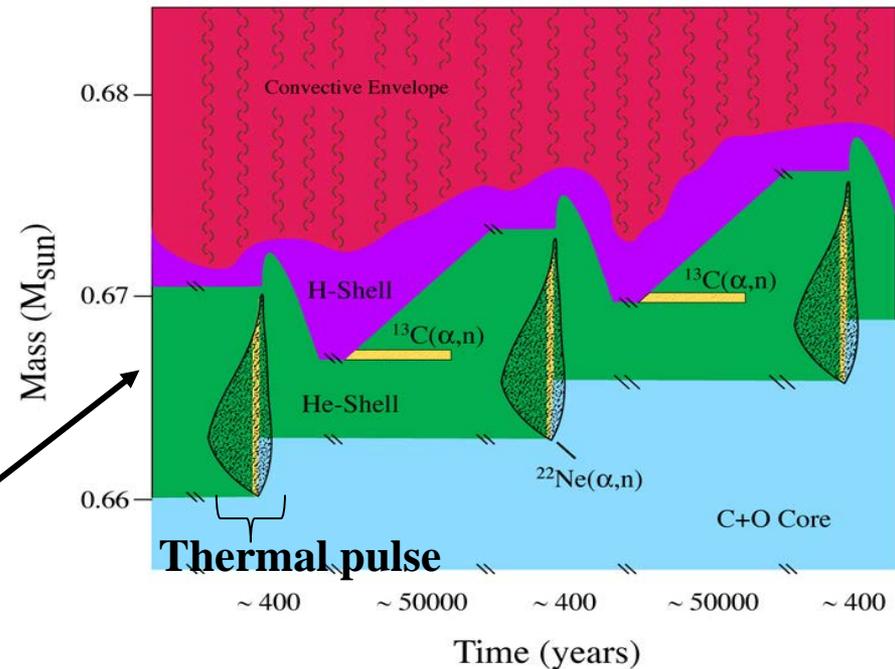
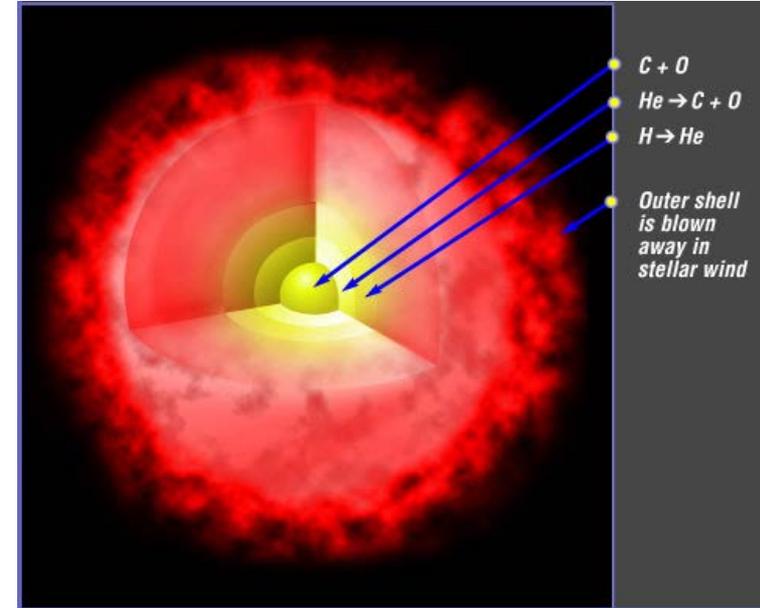
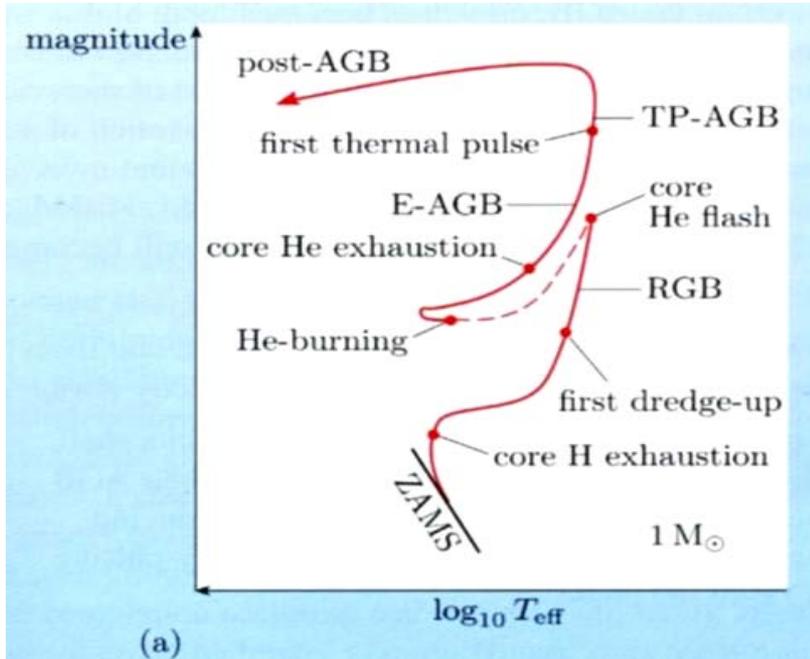
Stars of the horizontal branch (II)

- **Quiet ignition** of the He core (convective) for intermediate-mass stars ($2-10 M_{\odot}$)



Back to the HR diagram: AGB (Asymptotic Giant Branch) stars

- Asymptotic giant branch. E-AGB = early AGB
- H-burning ashes are brought to the surface by convection



Mixing of ashes from H and He burning.
Site of the s process (see later)

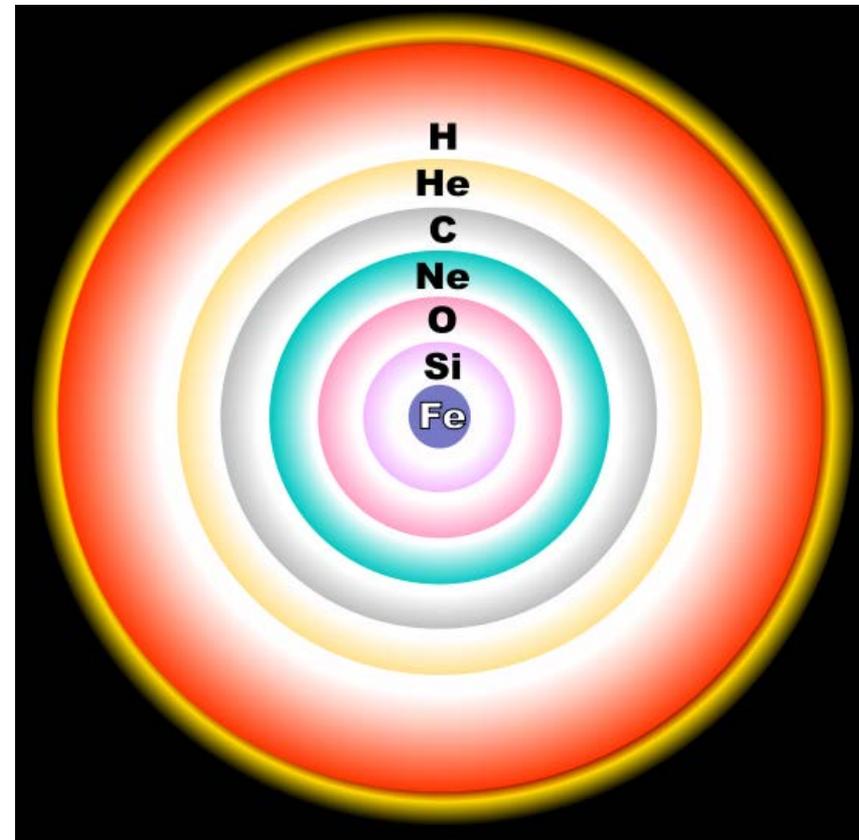
Evolution of a solar-type star

Time until the next stage (year)	T_C (MK)	T_{eff} (K)	ρ_C (g cm ⁻³)	Radius (R_{\odot})	Stellar stage
10^{10}	15	6000	10^2	1	Main sequence
10^8	50	4000	10^4	3	Subgiant
10^5	100	4000	10^5	50	Helium flash
5×10^7	200	5000	10^4	10	Horizontal branch
10^4	250	4000	10^5	200	AGB
10^5	300	100 000	10^7	0.01	Compact star enriched in C, O (planetary nebula)
–	100	50 000	10^7	0.01	White dwarf

Stellar Nucleosynthesis: Explosive stellar burning

Remember: “Onion Shell” Structure of Massive Star: Instant Before Core Collapse

- no other nuclear energy source is available to core
- core is in NSE, with $T=10^{10}$ K and $\rho=10^{10}$ g/cm³
- grows in mass; when it reaches 1.4 times solar mass, electron degeneracy pressure is unable to counteract gravity...
- rapid heating and compression of the matter enhance the collapse by 2 processes :
 - photodisintegration (many ways), exemple:
$$\gamma + {}^{56}\text{Fe} \rightarrow 13 {}^4\text{He} + 4\text{n}$$
 - electron capture on nuclei: ${}^A_Z\text{N} + e^- \rightarrow {}^A_{Z-1}\text{N} + \nu_e$
- when $\rho=10^{14}$ g/cm³: nuclei and nucleons feel short-range nuclear force [repulsive at very short distance] → inner part of core rebounds, producing an outward **moving shock wave...**



Explosive burning: → Shock moving through “Onion” Layers

Explosive Si burning:

➤ outgoing shock wave heats inner ^{28}Si layer of star to high T and ρ ; matter approaches NSE: composition entirely determined by values of T , ρ , and neutron excess η

- at $T=6$ GK, NSE is quickly established
- complete conversion of ^{28}Si to ^{56}Ni [“complete explosive silicon burning”]

Explosive O burning: ➤ next layer reached by shock is composed of ^{16}O

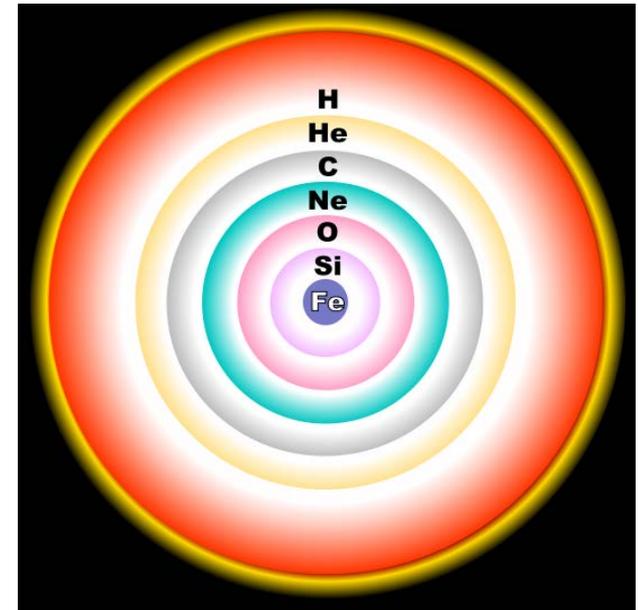
➤ at $T=3.8$ GK, ^{16}O fuel is depleted via $^{16}\text{O}+^{16}\text{O}$, $^{16}\text{O}(\gamma,\alpha)^{12}\text{C}$, etc., giving rise to two clusters in the mass regions of Si and Fe with more material in the Si region

➤ Most abundant nuclides: ^{28}Si , ^{32}S , ^{36}Ar , ^{40}Ca (“ α -element”) and some iron peak species

Explosive NeC burning: ➤ next layer reached by shock is composed of ^{16}O , ^{20}Ne , ^{12}C

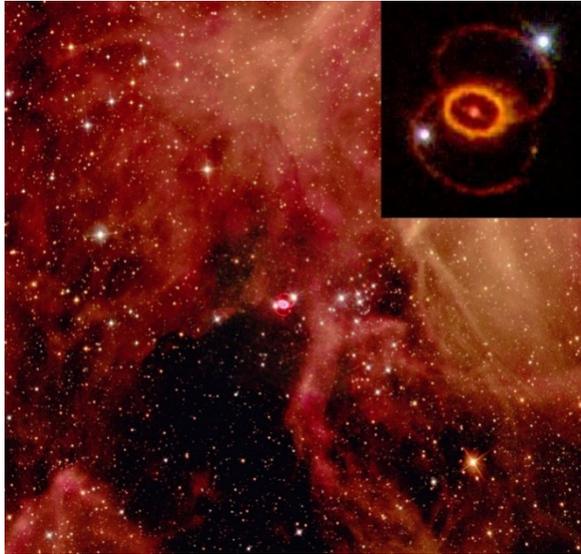
➤ at $T=2.5$ GK, ^{20}Ne & ^{12}C burn explosively

➤ Production of the radio-isotope ^{26}Al via $^{24}\text{Mg}(n,\gamma)^{25}\text{Mg}(p,\gamma)^{26}\text{Al}$ (γ -ray astronomy)

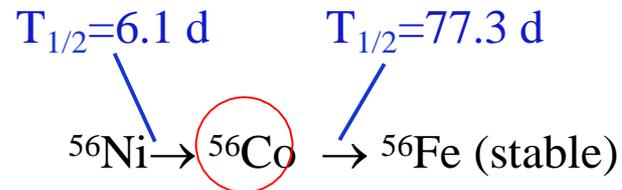
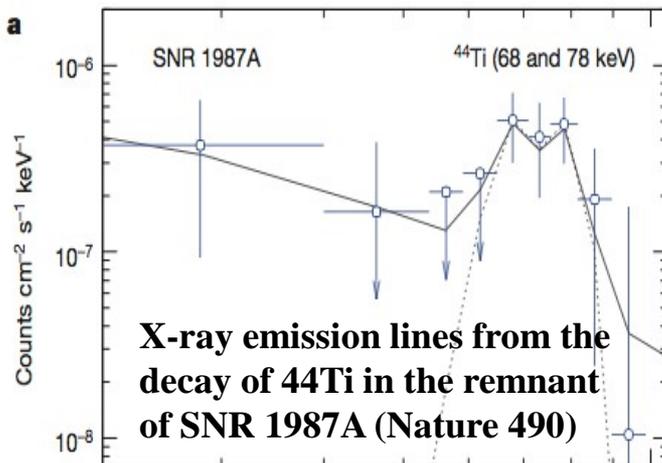
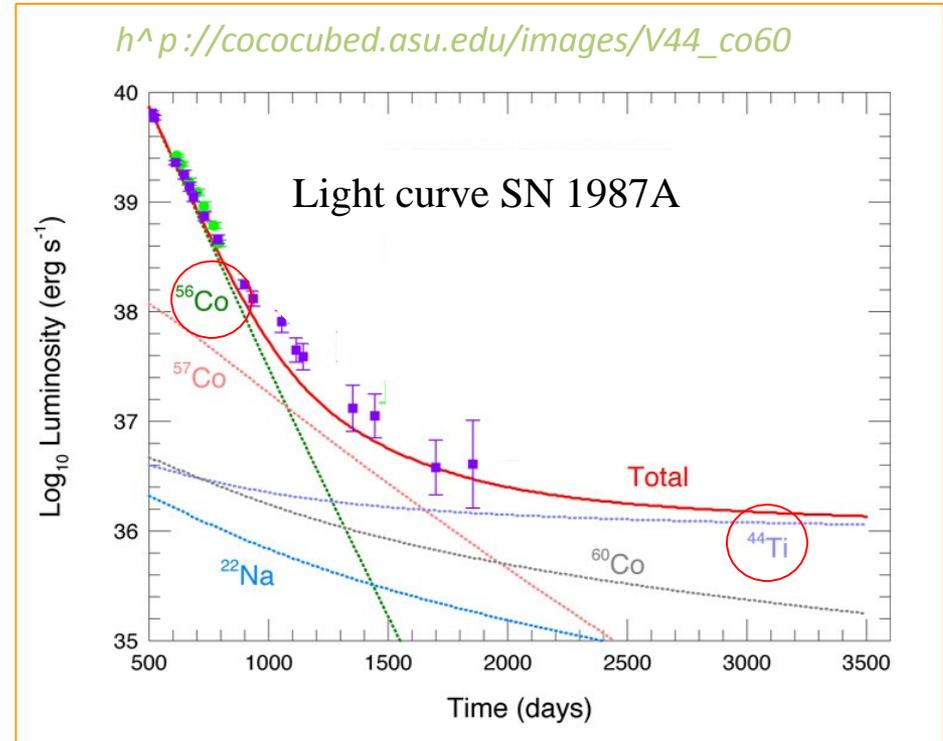


Explosive stellar burning: Core-Collapse Supernova observations

➤ The brightest objects in the universe → Luminosity : $\sim 10^{9-10} L_{\text{sun}}$



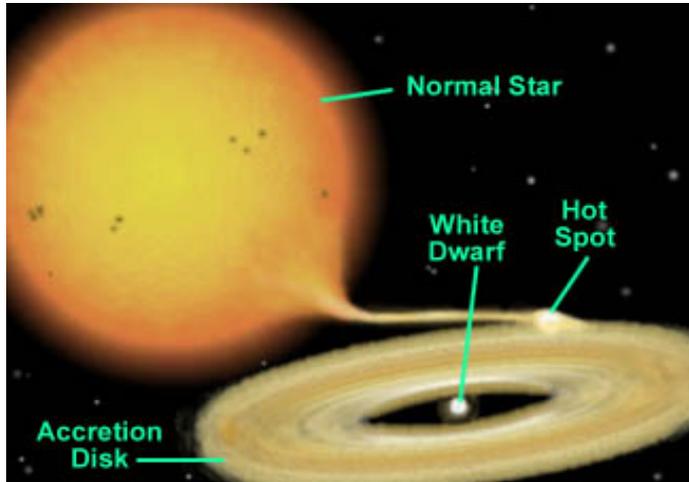
Supernova SN1987A



SN1987A: $m(^{56}\text{Ni}) = 0.07 \pm 0.01 M_{\odot}$

NOVAE = sudden increase in star's luminosity ($L \sim 10^4 - 10^6 L_{\odot}$ and $t \sim 1 \text{ h} - 1 \text{ d}$)

Binary system: White Dwarf + accreting normal star (e.g. SUN)



H-rich mass transfer from companion to WD



temperature and density increase
on WD's surface

thermonuclear runaway \Rightarrow cataclysmic explosion

(p, γ) reactions on proton-rich nuclei

$$T > 10^8 \text{ K}$$
$$\rho > 10^3 \text{ g cm}^{-3}$$

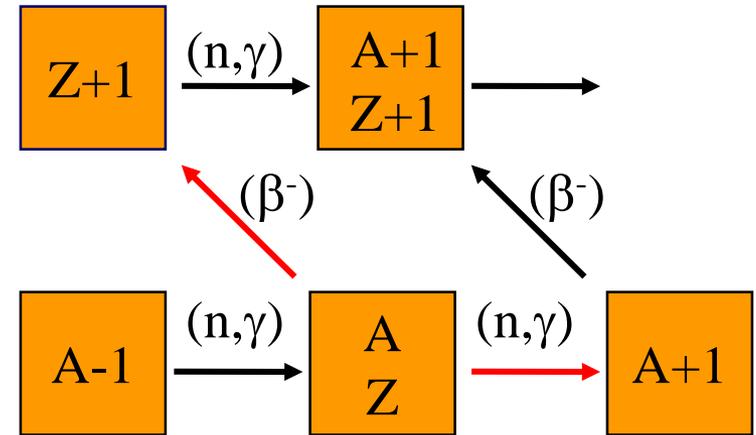


nucleosynthesis up to $A \sim 40$ mass region

Nucleosynthesis beyond Fe

Nucleosynthesis beyond Fe

- start with Fe *seeds*
- Elements heavier than Fe can not be synthesized by fusion reactions, their formation is explained by **neutrons capture reactions**
- Z change $\rightarrow \beta$ decay



DEFINE:

mean lifetime of nuclei A
against destruction by neutron capture

if $\tau_n > \tau_\beta \Rightarrow$ unstable nucleus decays
if $\tau_n < \tau_\beta \Rightarrow$ neutron capture

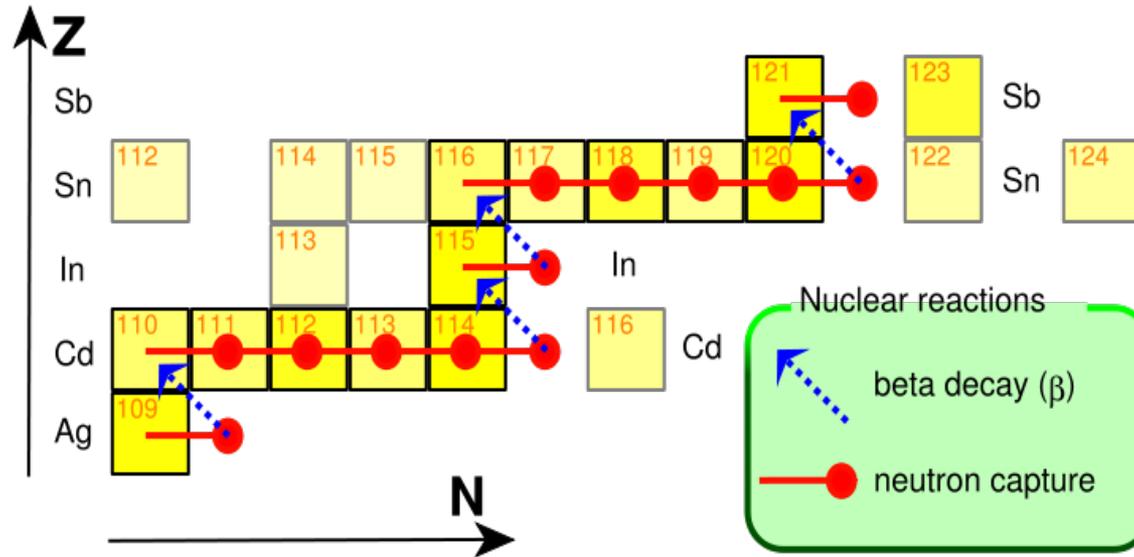
$$\tau_n(X) = \frac{1}{N_n \langle \sigma v \rangle}$$

τ_β = beta-decay lifetime
(seconds \rightarrow years)

s-process (s = slow neutron capture process)

unstable nucleus decays before capturing another neutron

$$\tau_{\beta} \ll \tau_n \Leftrightarrow N_n \sim 10^8 \text{ n/cm}^3$$

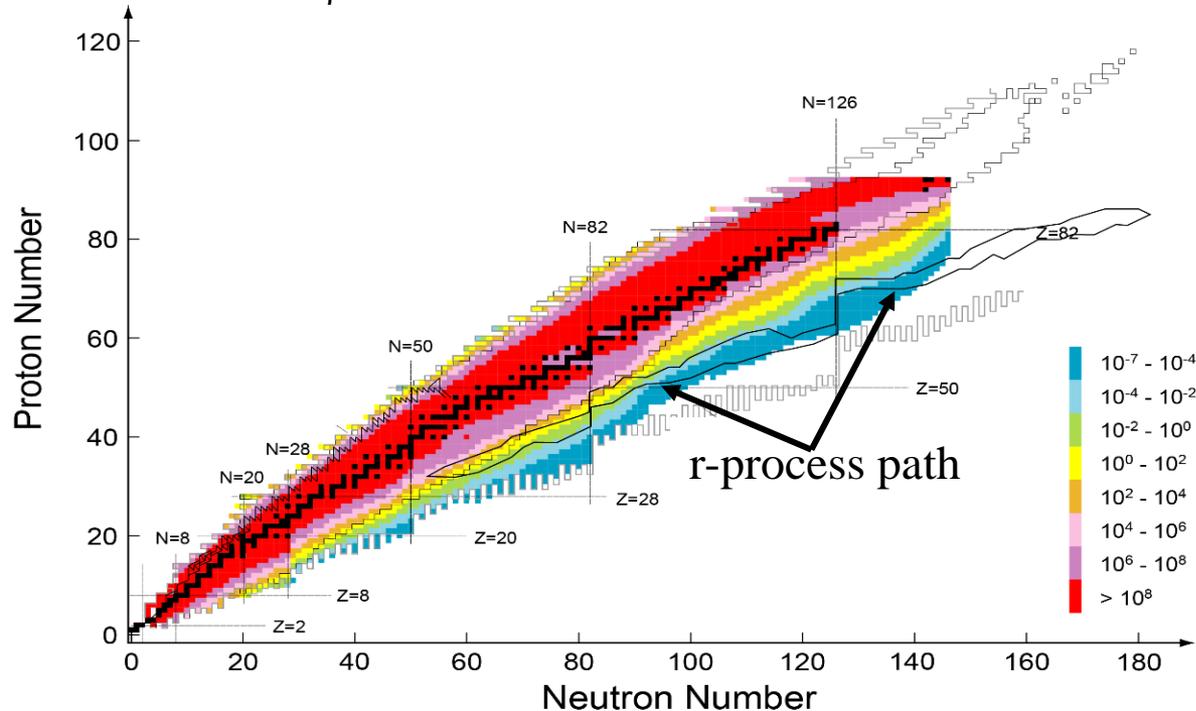


- synthesis path **along** valley of β - stability up to ^{209}Bi (long time scale $\sim 10^4$ years)
 - neutron source: $^{13}\text{C}(\alpha, n)^{16}\text{O}$ and/or $^{22}\text{Ne}(\alpha, n)^{25}\text{Mg}$
- **quiescent scenarios**: AGB stars (low and intermediate mass $M < 8 M_{\odot}$ or massive stars $M \geq 8 M_{\odot}$)

r-process (r = rapid neutron capture)

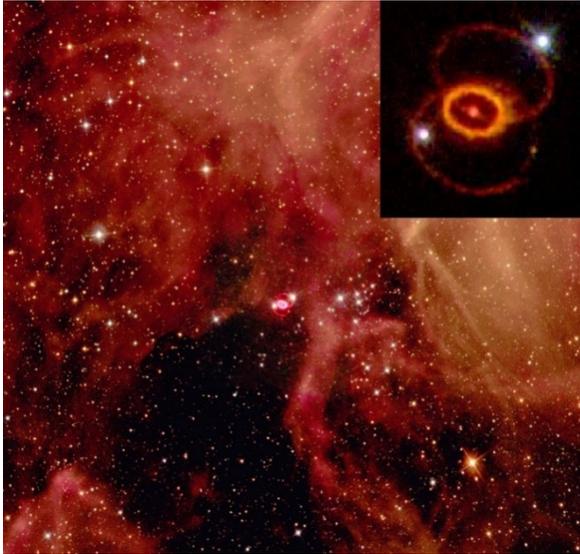
unstable nucleus captures another neutron before decaying

$$\tau_{\beta} \gg \tau_n \Leftrightarrow N_n \gg 10^{20} \text{ n/cm}^3$$



synthesis path far from valley of β –stability (short timescale \sim seconds)

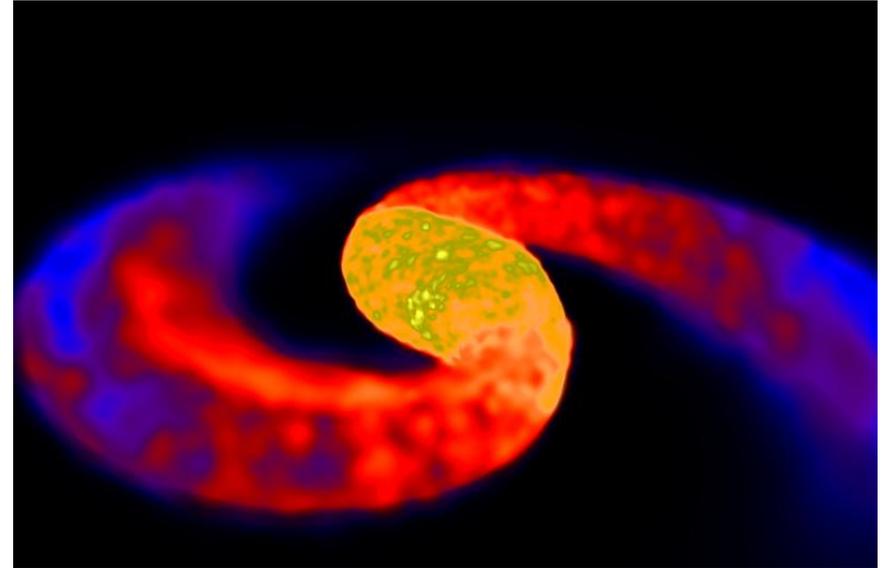
- synthesis of n-rich nuclei
- unknown nuclear properties
- explosive scenarios but where?



Supernova SN1987A

Core-collapse supernova?

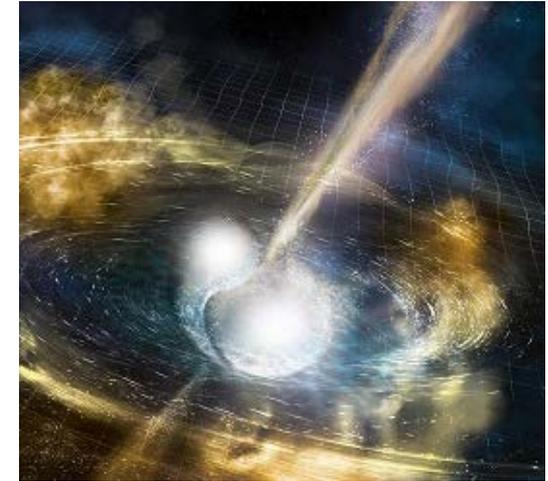
- Neutrino-winds from protoneutron stars, ...



Neutron star mergers ?

- Mergers are expected to eject around $0.01 M_{\odot}$ of very neutron rich-material
 - Sources of gravitational waves
- Observational signature: electromagnetic emission from radioactive decay of r-process nuclei
→ **kilonova**

- **17/08/2017** → gravitational waves were detected from neutron star mergers, **GW170817**, by LIGO/VIRGO.



- An optical transient source, Swope Supernova Survey 17a (SSS17a), was subsequently identified as the counterpart of this event in the galaxy NGC4993 at ~ 40 Mpc distance

First day

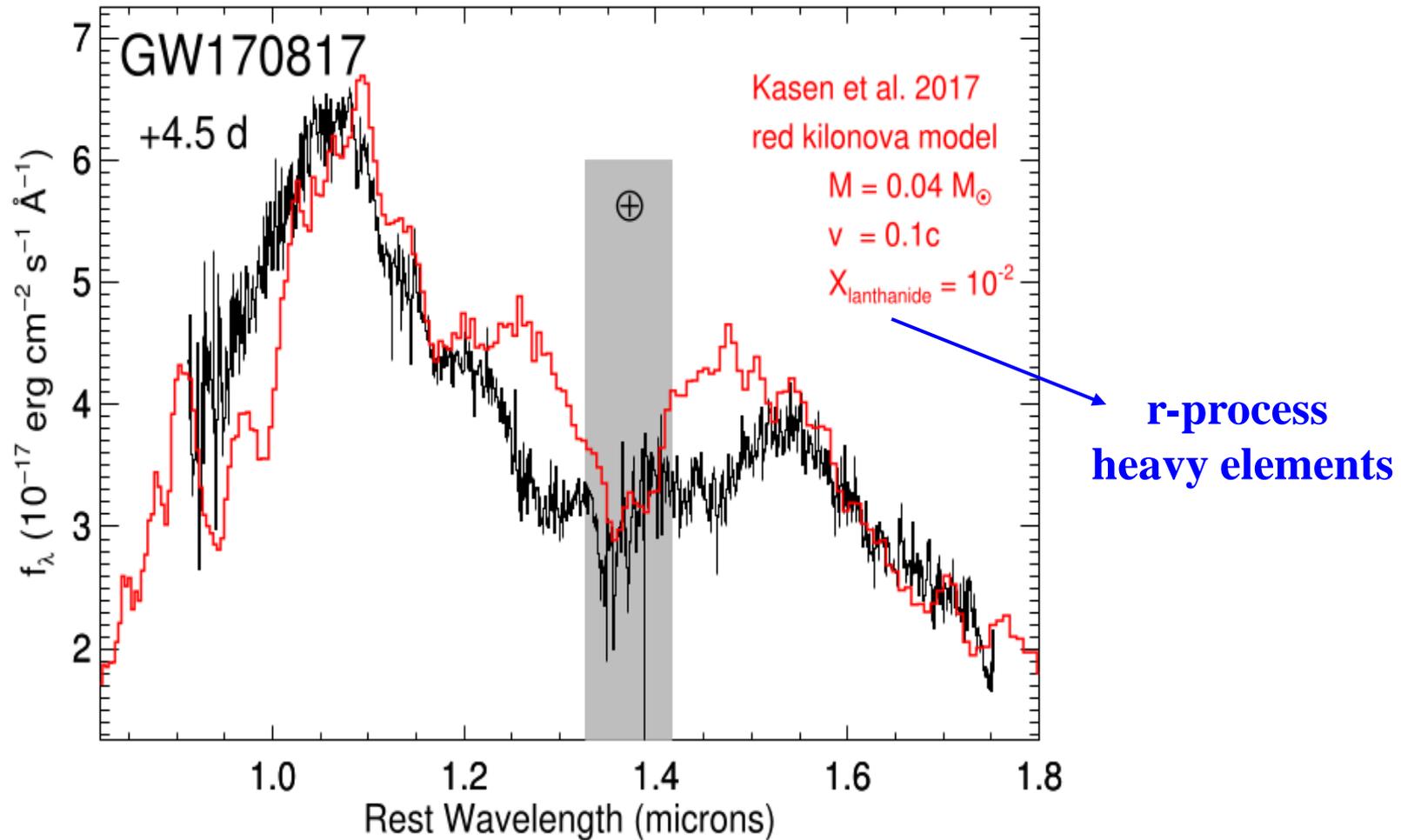
- Bright & Blue

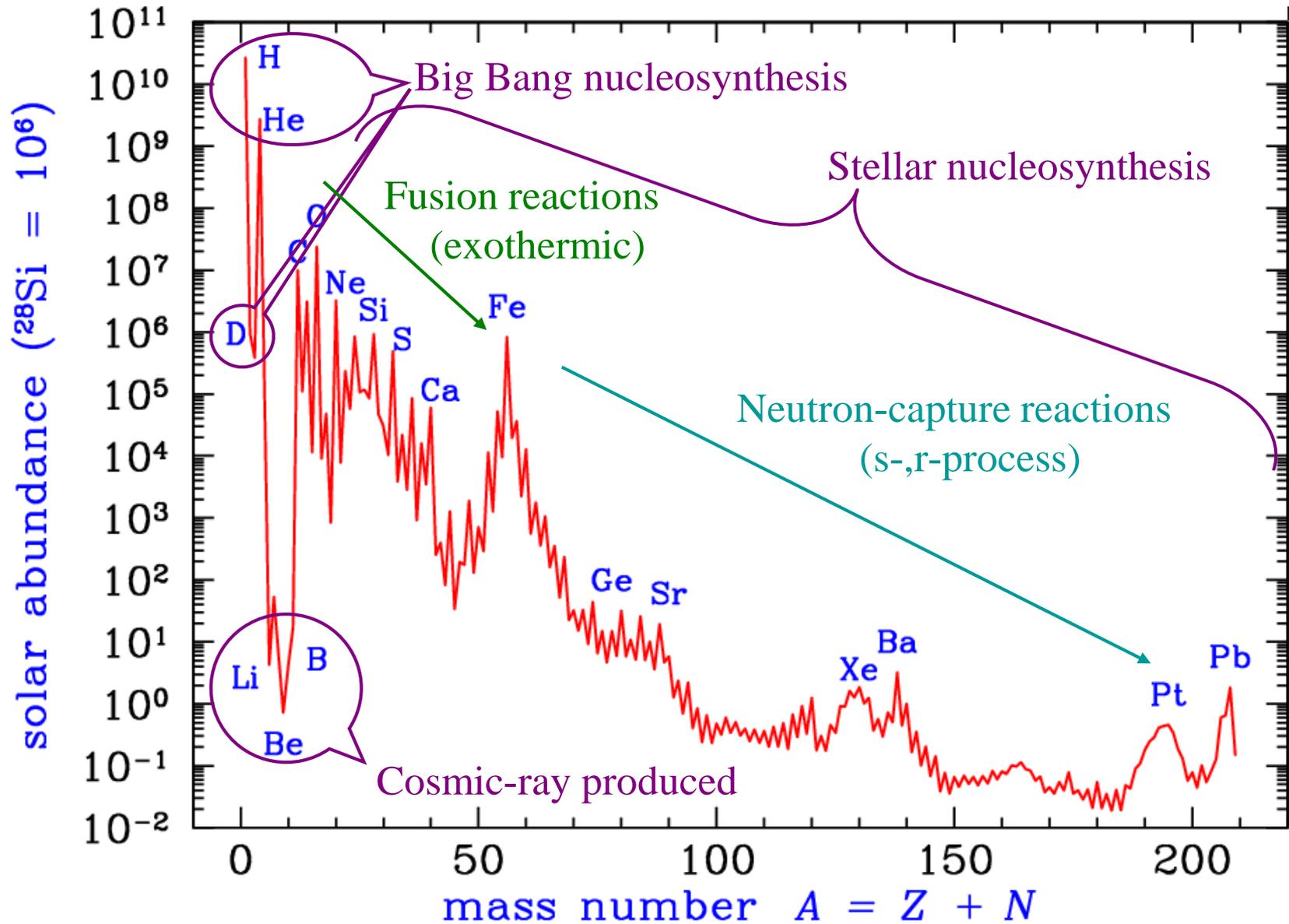
4 days later

- Fade & red

These observations show that the **heaviest r-process elements** were produced in neutron star mergers

Comparison of the measured near-infrared spectrum counterpart of the binary neutron star merger GW170817 versus red kilonova model





Summary:

Big-Bang & star life cycle

